

Automorphic string amplitudes

Henrik Gustafsson

HEP seminar

Université Libre de Bruxelles 2016

 hgustafsson.se

Based on

Small automorphic representations and degenerate Whittaker vectors

HG, Axel Kleinschmidt, Daniel Persson

[arXiv:1412.5625](#) [math.NT]

[GKP14]

Journal of Number Theory 166 (Sep, 2016) 344–399

Eisenstein series and automorphic representations

Philipp Fleig, HG, Axel Kleinschmidt, Daniel Persson

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Cambridge University Press (2017)

Upcoming work with

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$SL(n)$

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E_6, E_7, E_8

Outline

Compute Fourier coefficients of automorphic forms to capture information about non-perturbative effects such as instantons and black holes

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- ◎ Scattering amplitudes
4-graviton | Derivative expansion | U-duality | SUSY constraints

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Eisenstein series | Extracting physical information

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Parabolic subgroups | Limits of string theory | Adelic framework

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- Outlook

Motivation

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- Hecke eigenvalues
- Point counts of elliptic curves
- Langlands program
L-functions | The Langlands–Shahidi method

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- Statistical mechanics
Two-dimensional models of crystals

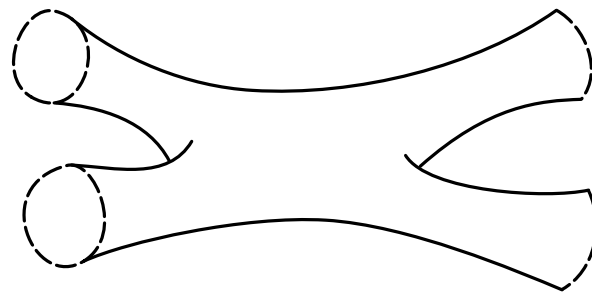
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String theory

Toroidal compactifications of type IIB string theory

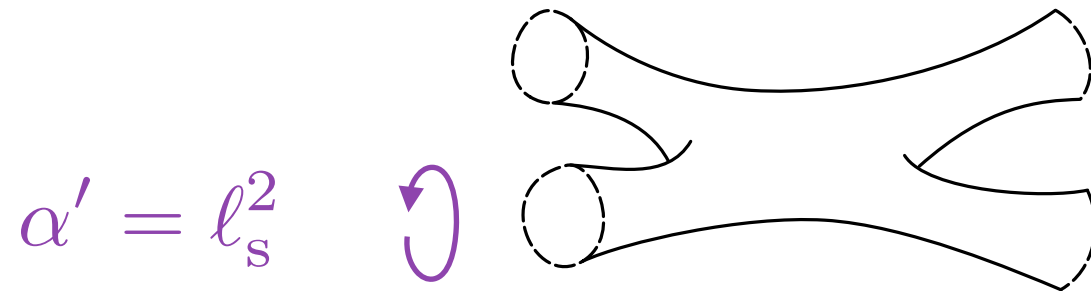
4-graviton scattering amplitudes



String theory

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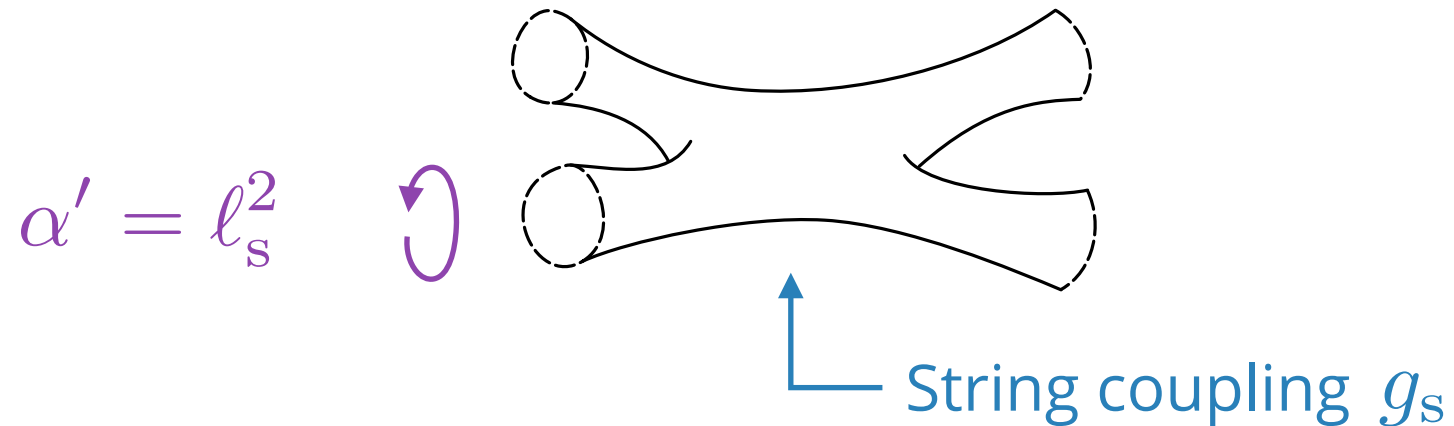
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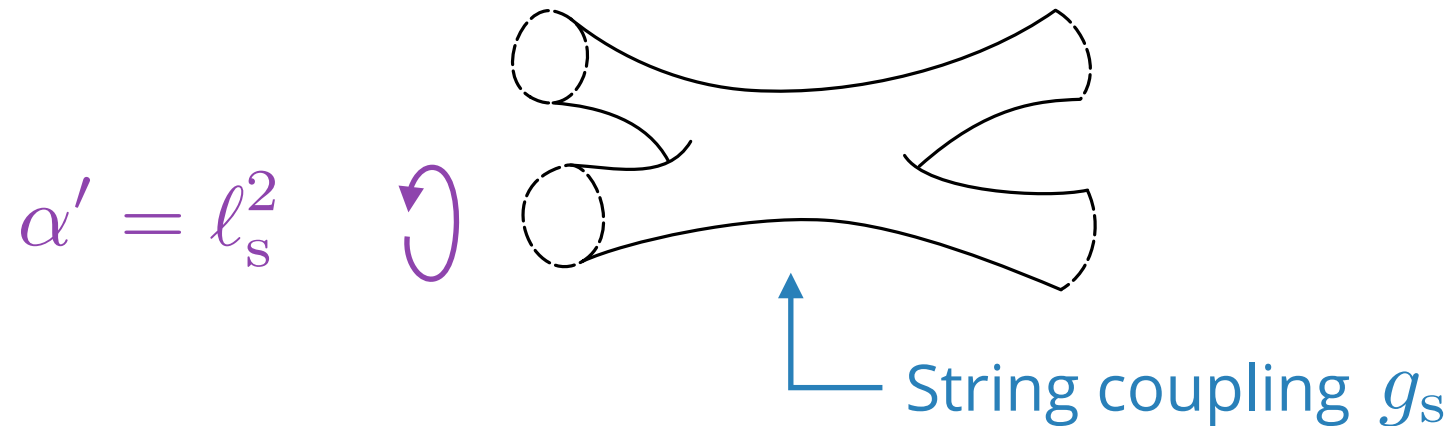
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String theory

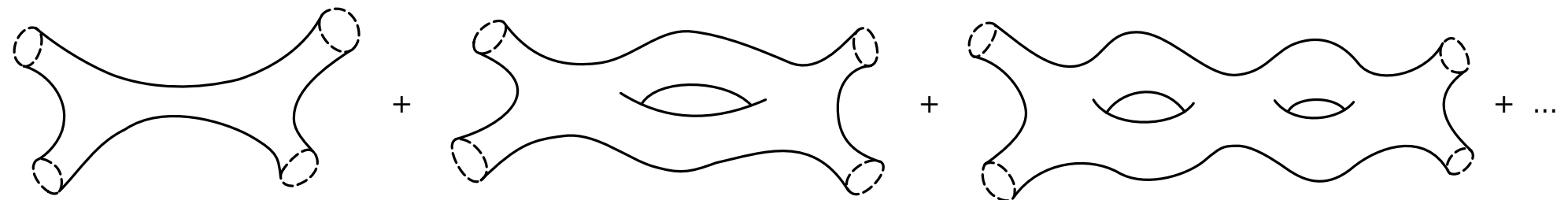
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4-graviton scattering amplitudes



$$s = -\frac{\alpha'}{4}(k_1 + k_2)^2 \quad t = -\frac{\alpha'}{4}(k_1 + k_3)^2 \quad u = -\frac{\alpha'}{4}(k_1 + k_4)^2$$

Interactions

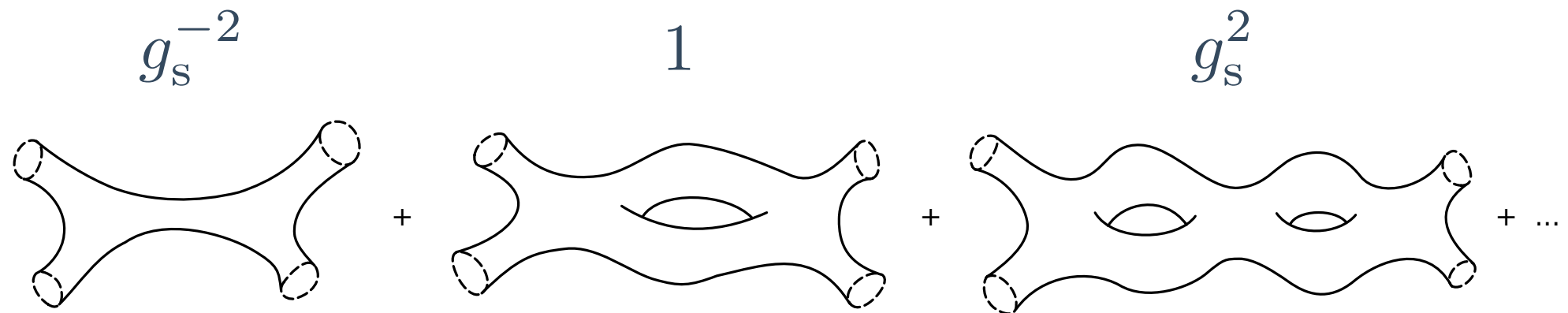
$$g_s^{-2} \quad 1 \quad g_s^2$$


The diagram shows a series of Feynman diagrams representing the 4-graviton amplitude in 10 dimensions. The first diagram is a tree-level exchange (s-channel) with a coefficient g_s^{-2} . The second diagram is a one-loop exchange (t-channel) with a coefficient 1. The third diagram is a two-loop exchange (t-channel) with a coefficient g_s^2 . The diagrams are separated by plus signs, and the sequence ends with an ellipsis.

4-graviton amplitude in 10 dimensions:

[Green-Schwarz, Green-Schwarz-Brink, Gross-Witten]

Interactions

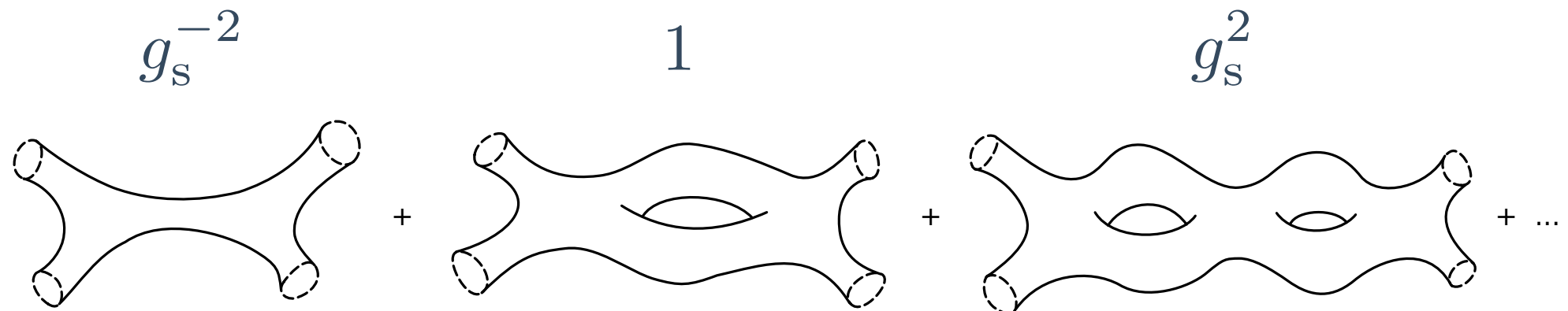


4-graviton amplitude in 10 dimensions:

$$\mathcal{A} = \left(g_s^{-2} \frac{1}{stu} \frac{\Gamma(1-s)\Gamma(1-t)\Gamma(1-u)}{\Gamma(1+s)\Gamma(1+t)\Gamma(1+u)} \right) \mathcal{R}^4$$

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Kinematic factor:
 Contraction of 4 linearized
 Riemann tensors
 $t_8 t_8 \mathcal{R}^4$

[Green-Schwarz, Green-Schwarz-Brink, Gross-Witten]

Interactions

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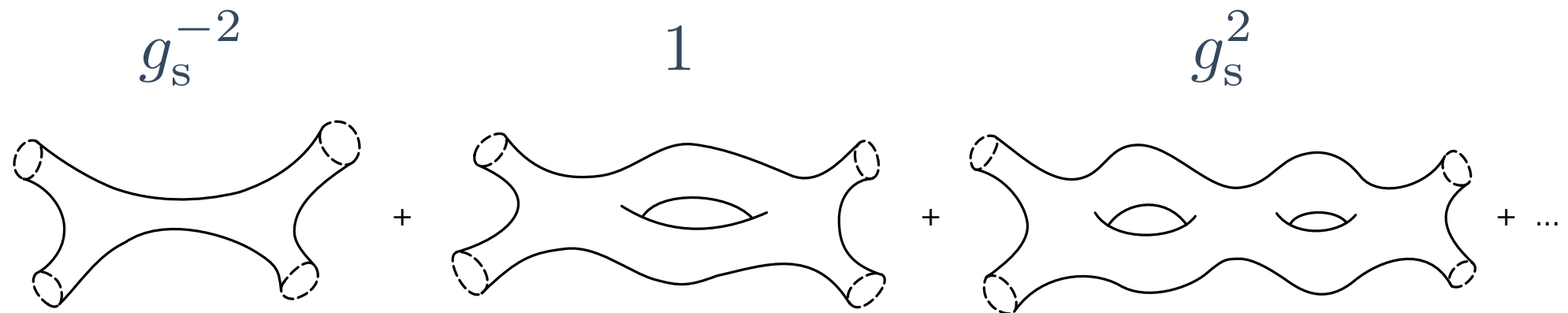
The diagram shows a series of Feynman diagrams representing 4-graviton interactions in string theory. The first diagram is a sphere with four external legs, labeled g_s^{-2} . The second diagram is a torus with four external legs, labeled 1. The third diagram is a genus-2 surface (two holes) with four external legs, labeled g_s^2 . The diagrams are summed together with ellipses indicating higher-order terms.

4-graviton amplitude in 10 dimensions:

$$\mathcal{A} = \left(g_s^{-2} \frac{1}{stu} \frac{\Gamma(1-s)\Gamma(1-t)\Gamma(1-u)}{\Gamma(1+s)\Gamma(1+t)\Gamma(1+u)} + 2\pi \int_{\mathcal{F}} \frac{d^2\tau}{(\text{Im } \tau)^2} \mathcal{B}_1(s, t, u; \tau) \right) \mathcal{R}^4$$

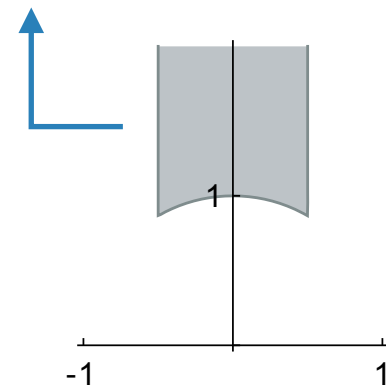
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Interactions



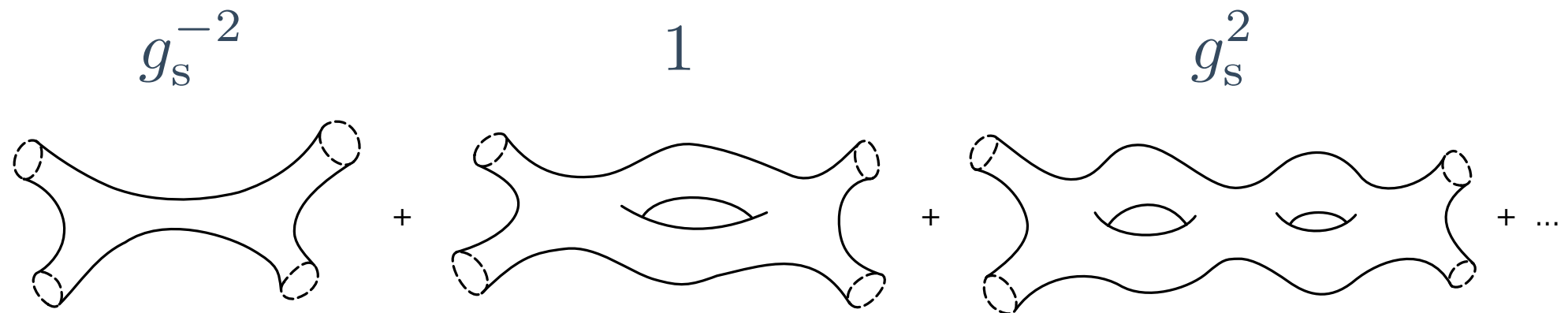
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Interactions



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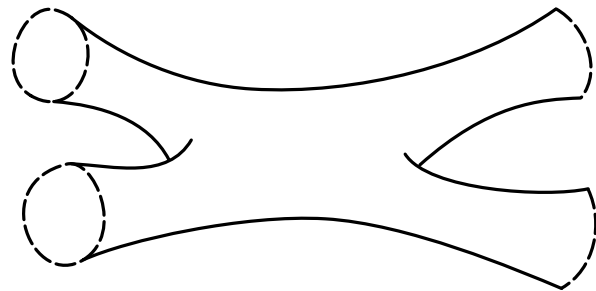
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The diagram shows the fundamental domain $\mathcal{F}$ in the complex plane, which is a shaded gray region bounded by $\tau = -1$, $\tau = 1$, and $\tau = i$. A vertical line segment is drawn from the origin to the top boundary at $\tau = i$, with a tick mark at 1. An arrow points from the label $\mathcal{F}$ to the domain, and another arrow points from the label "Modular invariant function" to the integrand $\mathcal{B}_1(s, t, u; \tau)$.

[Green-Schwarz, Green-Schwarz-Brink, Gross-Witten]

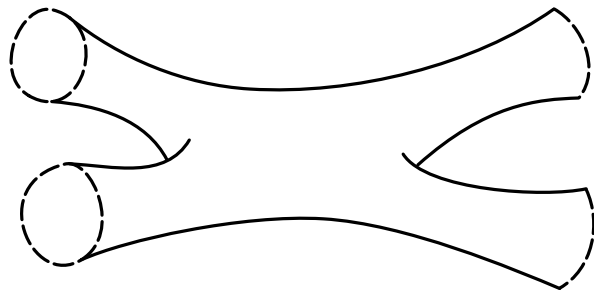
Interactions

String theory



Interactions

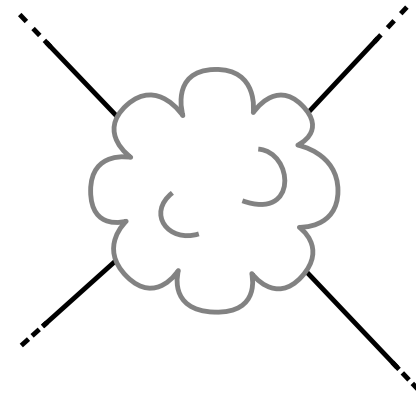
String theory



Effective field theory

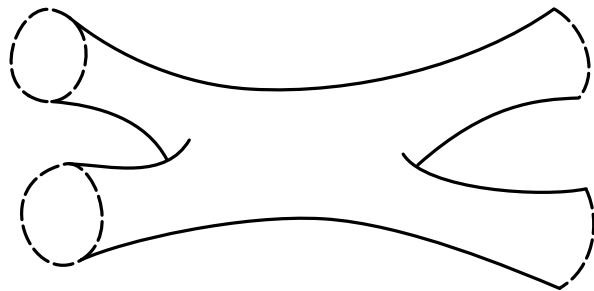


Supergravity + $\mathcal{O}(\alpha')$



Interactions

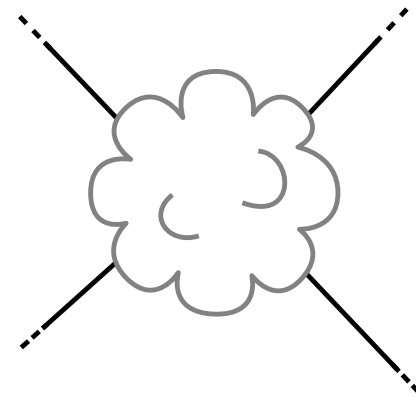
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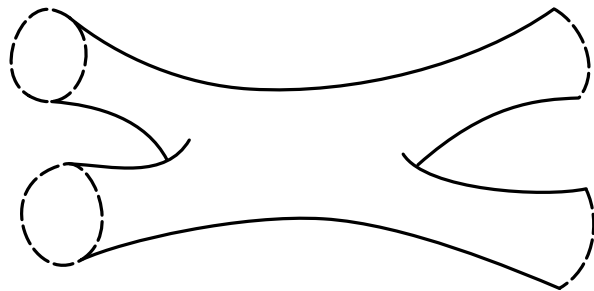
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$$s, t, u = \mathcal{O}(\alpha') \quad p \longleftrightarrow \partial \implies \alpha' \text{-expansion} = \partial \text{-expansion}$$

Interactions

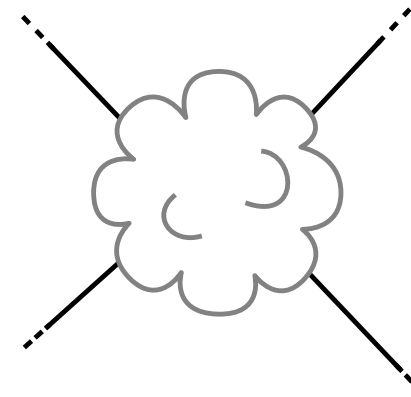
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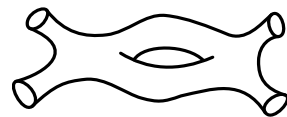
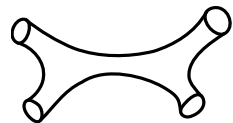
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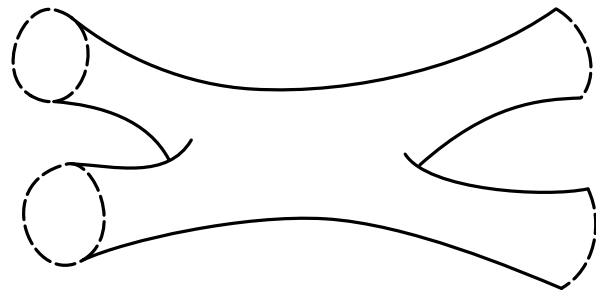


(Einstein frame)

$$\mathcal{L} \propto R + (\alpha')^3 \left(2\zeta(3)g_s^{-3/2} + 4\zeta(2)g_s^{1/2} + \dots \right) R^4 + \dots$$

Interactions

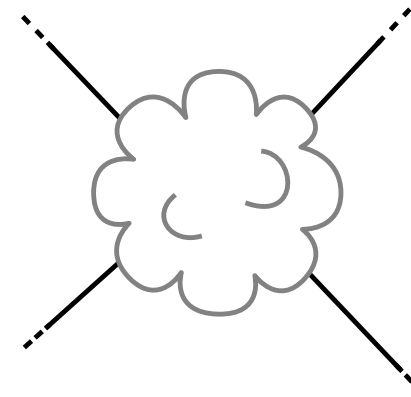
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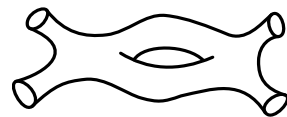
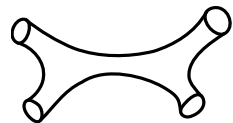
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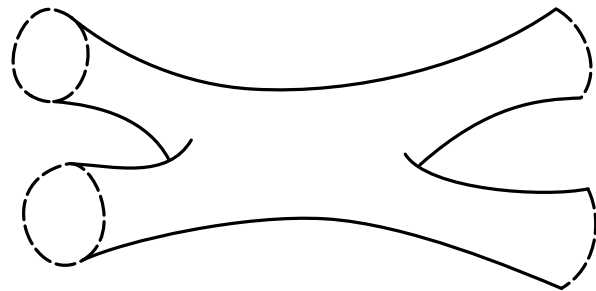
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Contraction of 4 Riemann tensors



Interactions

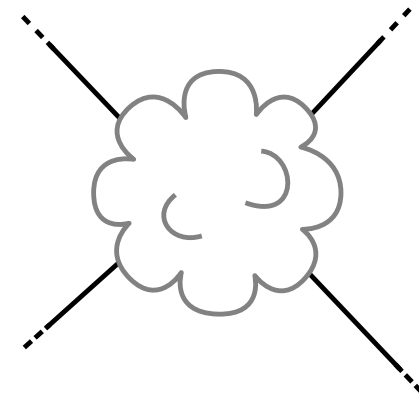
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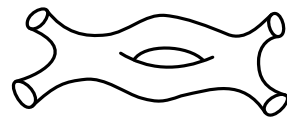
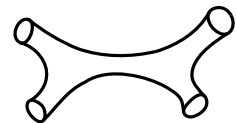
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(Einstein frame)

$$\begin{aligned} \mathcal{L} \propto R &+ (\alpha')^3 \left(2\zeta(3)g_s^{-3/2} + 4\zeta(2)g_s^{1/2} + \dots \right) R^4 + \\ &(\alpha')^5 \left(\zeta(5)g_s^{-5/2} + \dots \right) D^4 R^4 + \\ &(\alpha')^6 \left(\frac{2}{3}\zeta(3)^2 g_s^{-3} + \frac{4}{3}\zeta(2)\zeta(3)g_s^{-1} + \dots \right) D^6 R^4 + \mathcal{O}((\alpha')^7) \end{aligned}$$

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$$\mathcal{L} \propto R + (\alpha')^3 \mathcal{E}_0(\tau) R^4 + (\alpha')^5 \mathcal{E}_4(\tau) D^4 R^4 + (\alpha')^6 \mathcal{E}_6(\tau) D^6 R^4 + \dots$$

$$\tau = \chi + i g_s^{-1}$$

Interactions

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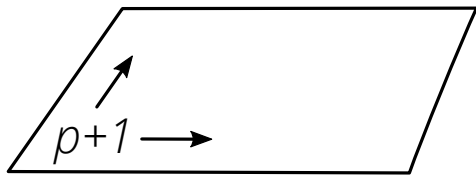
$$\tau = \chi + i g_s^{-1}$$

Bao, Basu, Bossard, Cederwall, Fleig, Green, Gubay, Gutperle, HG, Kazhdan, Kiritsis, Kleinschmidt, Lambert, Miller, Nilsson, Obers, Persson, Pioline, Russo, Sethi, Vanhove, Verschinin, Waldron, West, ...

Non-perturbative effects

[Green, Polchinski]

Non-perturbative effects

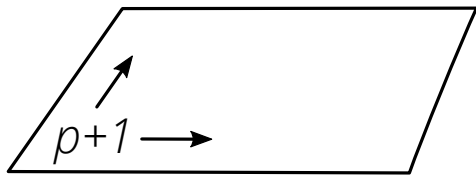


Dp -brane

p space directions

1 time direction

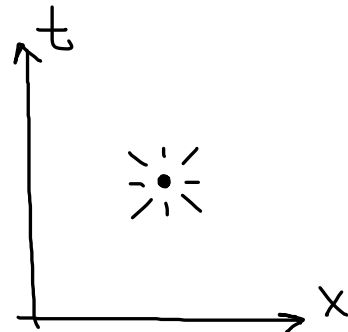
Non-perturbative effects



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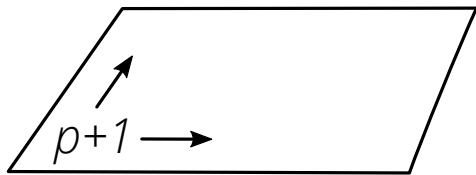


D-instanton

$p = -1$

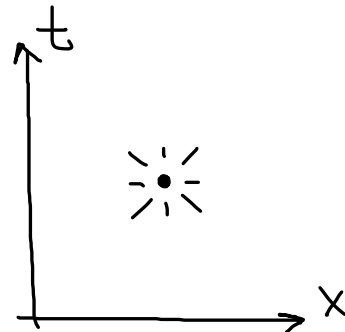
[Green, Polchinski]

Non-perturbative effects



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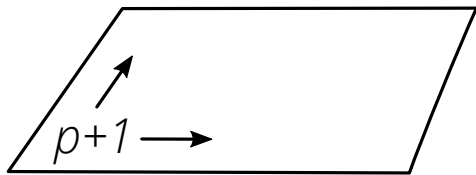
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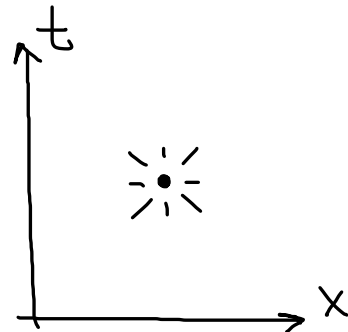
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Non-perturbative effects



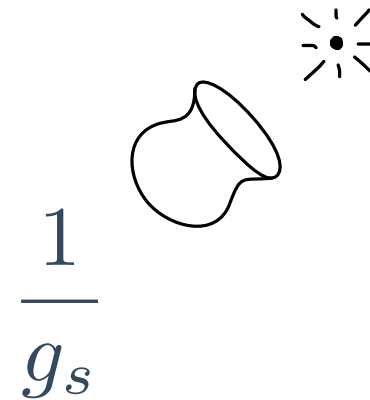
Dp-brane

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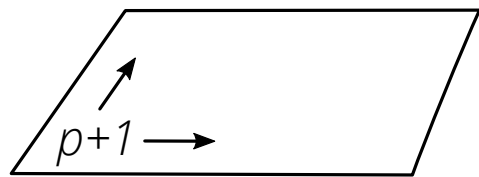
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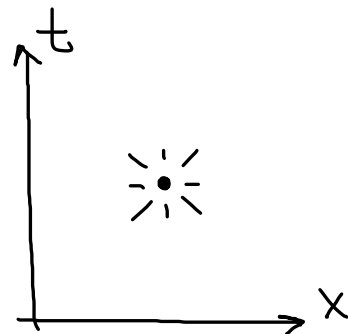
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Non-perturbative effects



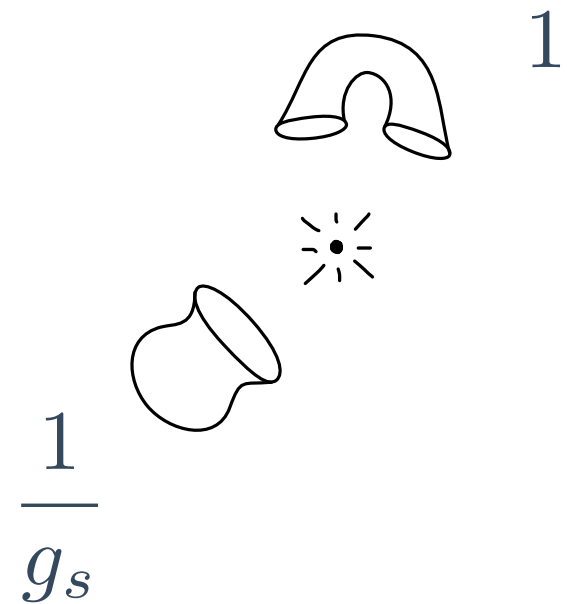
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p space directions
1 time direction



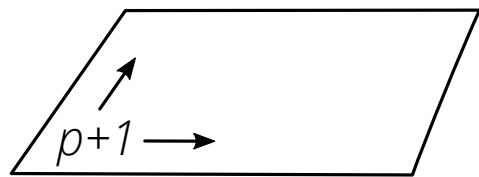
D-instanton

$p = -1$



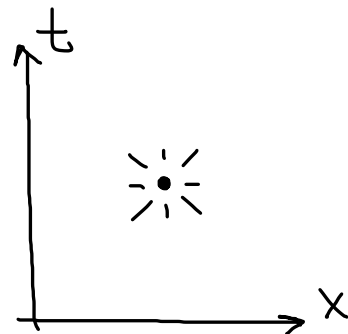
[Green, Polchinski]

Non-perturbative effects



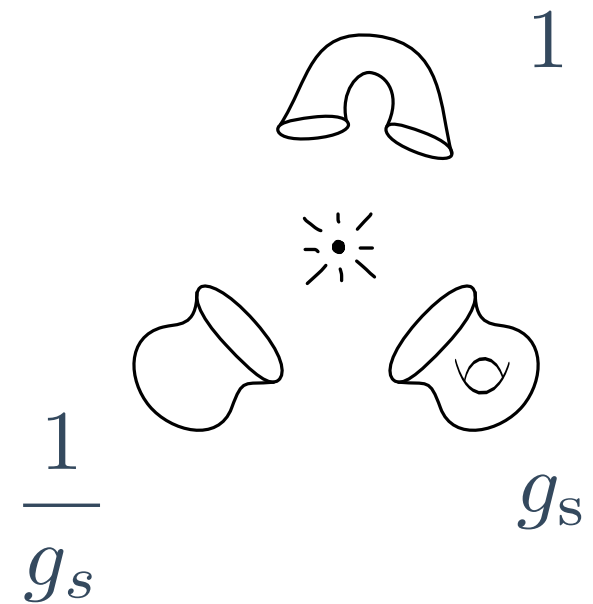
Dp-brane

p space directions
1 time direction



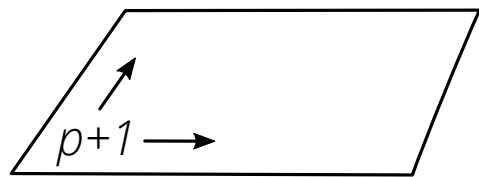
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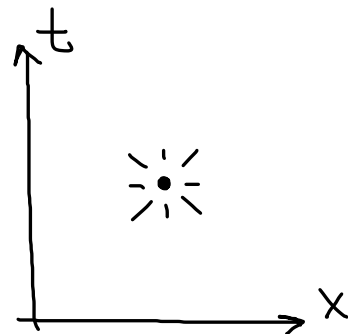
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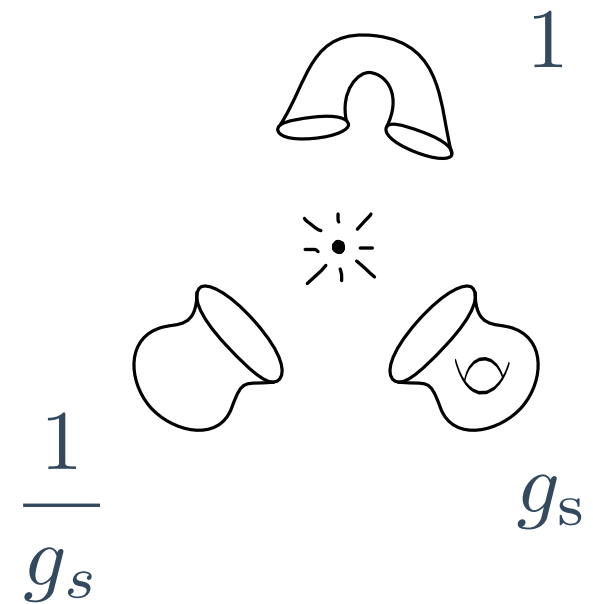
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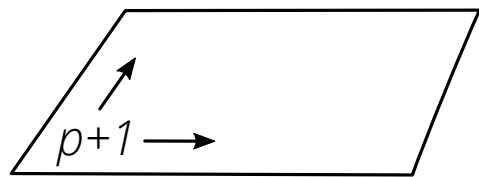
$p = -1$



1

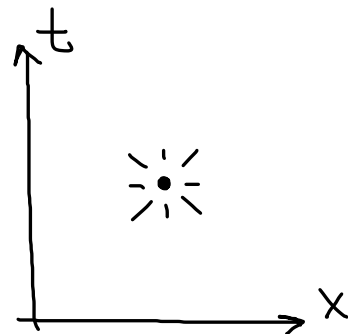
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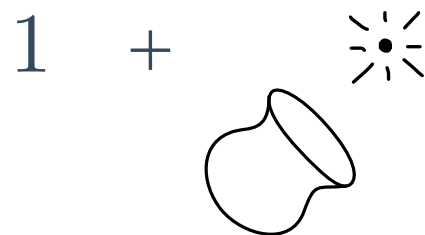
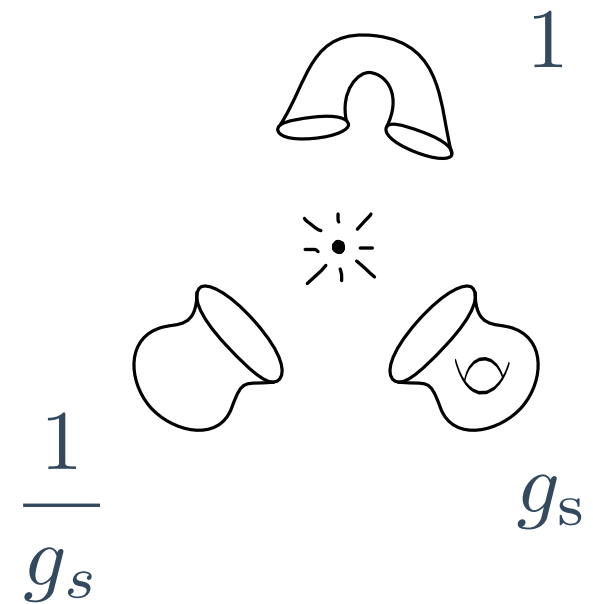
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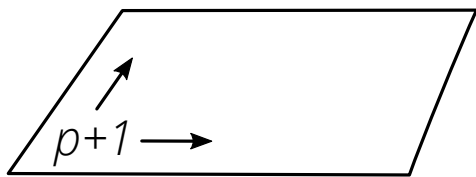
D-instanton

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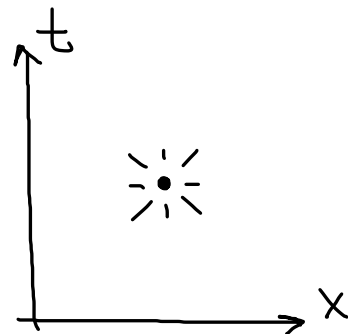
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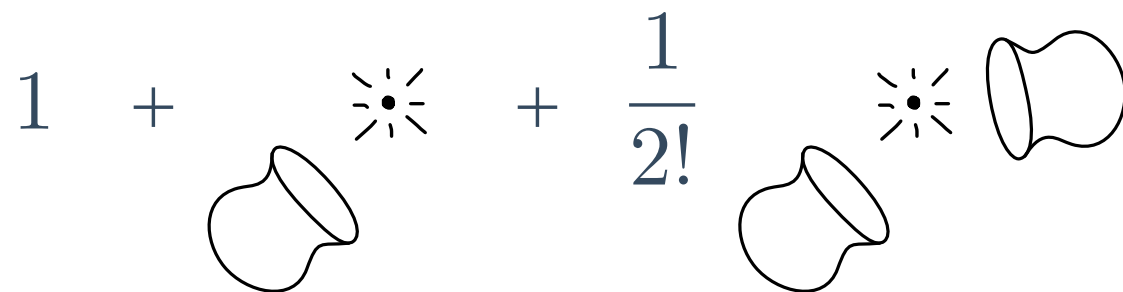
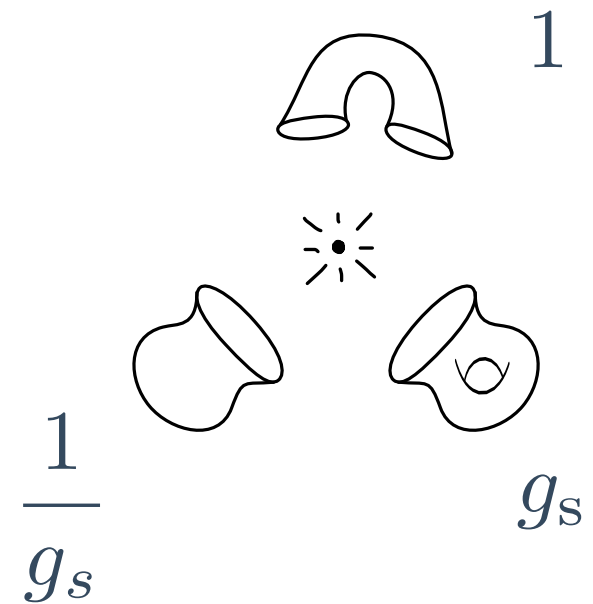
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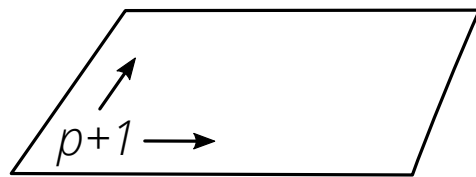
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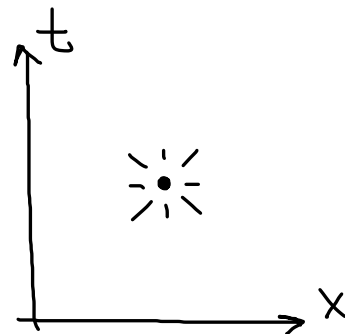
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Non-perturbative effects



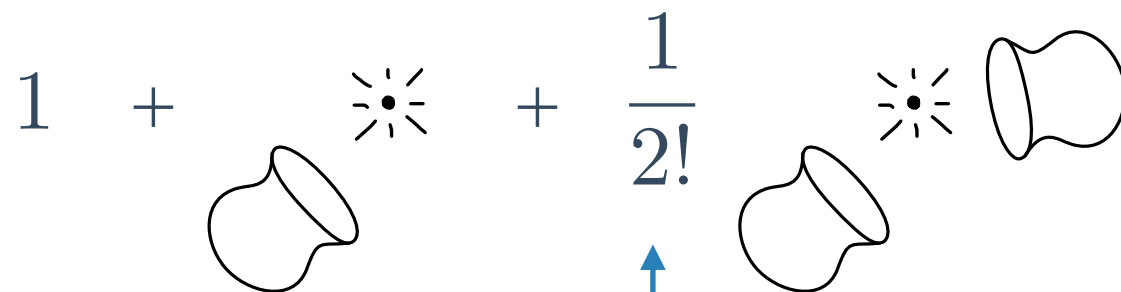
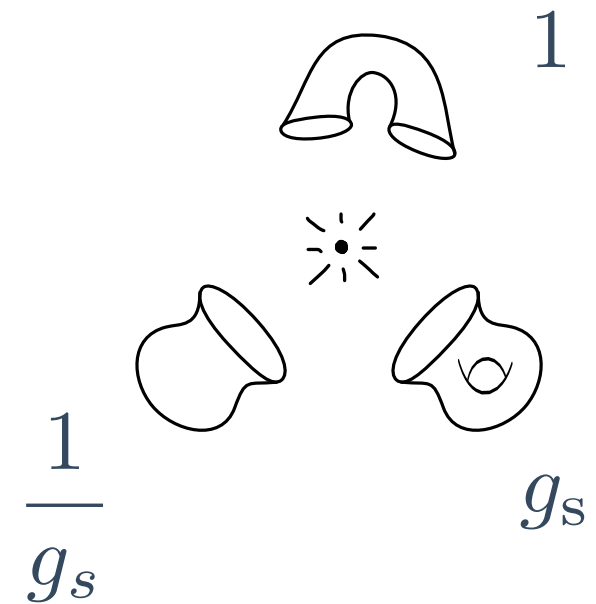
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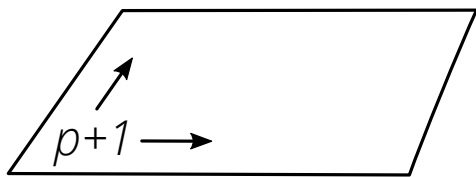
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Symmetry factor for identical disks

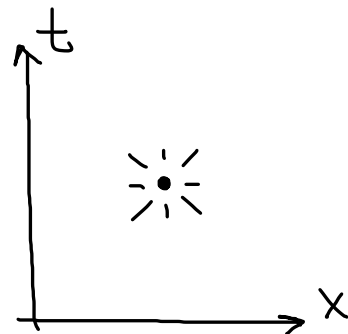
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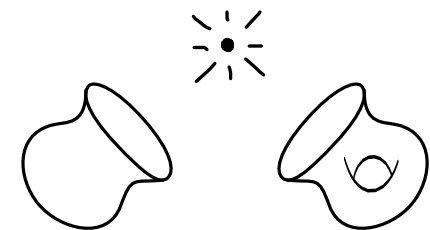
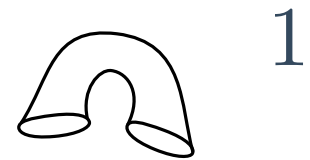
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1 time direction



D-instanton

$p = -1$



$\frac{1}{g_s}$

$$1 + \text{disk} + \frac{1}{2!} \text{disk} \times \text{disk} + \frac{1}{3!} \text{disk} \times \text{disk} \times \text{disk} + \dots$$

↑ Symmetry factor for identical disks

[Green, Polchinski]

Non-perturbative effects

$$1 + \text{[diagram 1]} + \frac{1}{2!} \text{[diagram 2]} + \frac{1}{3!} \text{[diagram 3]} + \dots$$

The diagram shows a series of terms in a mathematical expansion. The first term is 1. The second term is a plus sign followed by a diagram of a single torus (a donut shape) with a central point and eight radiating lines, representing a singularity. The third term is a plus sign followed by the fraction 1/2! followed by a diagram of two tori connected by a vertical line, with a central point and eight radiating lines. The fourth term is a plus sign followed by the fraction 1/3! followed by a diagram of three tori arranged in a triangle, with a central point and eight radiating lines. The series ends with an ellipsis (...).

[Green-Gutperle]

Non-perturbative effects

$$1 + \text{[diagram 1]} + \frac{1}{2!} \text{[diagram 2]} + \frac{1}{3!} \text{[diagram 3]} + \dots$$

The diagrams represent instanton configurations:

- Diagram 1: A single instanton (a sphere with a cap).
- Diagram 2: Two instantons connected by a line with a central dot and radiating lines, representing a two-instanton configuration.
- Diagram 3: Three instantons arranged in a triangle, connected by lines with central dots and radiating lines, representing a three-instanton configuration.

$$\exp \left(\text{[diagram 1]} \right)$$

Non-perturbative effects

$$1 + \text{[torus]} \star + \frac{1}{2!} \text{[torus]} \star \text{[torus]} + \frac{1}{3!} \text{[torus]} \star \text{[torus]} \star \text{[torus]} + \dots$$

$$\exp \left(\text{[torus]} \right) \sim \exp \left(-\frac{\text{const}}{g_s} \right)$$

Non-perturbative effects

$$1 + \text{[one-loop diagram]} + \frac{1}{2!} \text{[two-loop diagram]} + \frac{1}{3!} \text{[three-loop diagram]} + \dots$$

$$\exp\left(\text{[one-loop diagram]}\right) \sim \exp\left(-\frac{\text{const}}{g_s}\right) \quad \text{Non-perturbative in } g_s$$

Non-perturbative effects

$$1 + \text{[one-loop diagram]} \star + \frac{1}{2!} \text{[one-loop diagram]} \star \text{[two-loop diagram]} + \frac{1}{3!} \text{[one-loop diagram]} \star \text{[two-loop diagram]} + \dots$$

$$\exp\left(\text{[one-loop diagram]}\right) \sim \exp\left(-\frac{\text{const}}{g_s}\right) \quad \text{Non-perturbative in } g_s$$

$$\mathcal{E}_0(\tau) = 2\zeta(3)\tau_2^{3/2} + 4\zeta(2)\tau_2^{-1/2} + \dots + \underbrace{Ce^{2\pi i\tau}}_{\dots\dots\dots} + \dots$$

$$\tau = \tau_1 + i\tau_2 = \chi + ig_s^{-1}$$

[Green-Gutperle]

Moduli space

$$R + (\alpha')^3 \mathcal{E}_0^{(D)}(g) R^4 + (\alpha')^5 \mathcal{E}_4^{(D)}(g) D^4 R^4 + (\alpha')^6 \mathcal{E}_6^{(D)}(g) D^6 R^4 + \dots$$

$$\mathcal{M}_{\text{classical}} = G(\mathbb{R})/K$$

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$$\mathcal{M}_{\text{classical}} = G(\mathbb{R})/K$$

D	$G(\mathbb{R})$	K
10	$SL(2, \mathbb{R})$	$SO(2)$
9	$SL(2, \mathbb{R}) \times \mathbb{R}^+$	$SO(2)$
8	$SL(3, \mathbb{R}) \times SL(2, \mathbb{R})$	$SO(3) \times SO(2)$
7	$SL(5, \mathbb{R})$	$SO(5)$
6	$Spin(5, 5; \mathbb{R})$	$(Spin(5) \times Spin(5))/\mathbb{Z}_2$
5	$E_6(\mathbb{R})$	$USp(8)/\mathbb{Z}_2$
4	$E_7(\mathbb{R})$	$SU(8)/\mathbb{Z}_2$
3	$E_8(\mathbb{R})$	$Spin(16)/\mathbb{Z}_2$

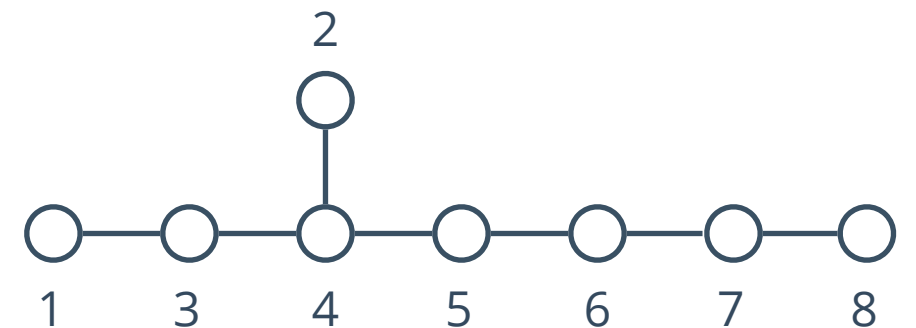
[Cremmer-Julia]

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[Cremmer-Julia]

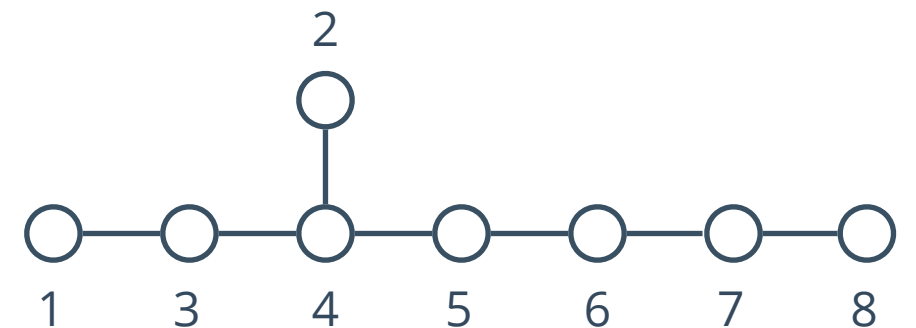
Moduli space

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$$\mathcal{M}_{\text{classical}} = G(\mathbb{R})/K$$

$$\mathcal{E}_n(\tau) = \mathcal{E}_n^{(10)}(g)$$

D	$G(\mathbb{R})$	K
10	$SL(2, \mathbb{R})$	$SO(2)$
9	$SL(2, \mathbb{R}) \times \mathbb{R}^+$	$SO(2)$
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Moduli space

10 dimensions:

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$$\tau = \chi + ig_s^{-1} \in \mathbb{H} = \{z \in \mathbb{C} \mid \text{Im } z > 0\} \cong SL(2, \mathbb{R})/SO(2, \mathbb{R})$$

Moduli space

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$$\tau = \chi + ig_s^{-1} \in \mathbb{H} = \{z \in \mathbb{C} \mid \text{Im } z > 0\} \cong SL(2, \mathbb{R})/SO(2, \mathbb{R})$$

No similar structure for lower dimensions

U-duality

$G(\mathbb{R}) \curvearrowright \mathcal{M}_{\text{classical}}$ classical symmetry

[Hull-Townsend]

U-duality

$G(\mathbb{R}) \curvearrowright \mathcal{M}_{\text{classical}}$ classical symmetry

Quantization of charges

[Hull-Townsend]

U-duality

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Quantization of charges $\implies$ classical symmetry $\longrightarrow$ discrete symmetry

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$G(\mathbb{R})$

Chevalley group $G(\mathbb{Z})$

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7	$SL(5, \mathbb{R})$	$SO(5)$	$SL(5, \mathbb{Z})$
6	$Spin(5, 5; \mathbb{R})$	$(Spin(5) \times Spin(5)) / \mathbb{Z}_2$	$Spin(5, 5; \mathbb{Z})$
5	$E_6(\mathbb{R})$	$USp(8) / \mathbb{Z}_2$	$E_6(\mathbb{Z})$
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All observables are invariant under $G(\mathbb{Z})$

[Hull-Townsend]

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$$\mathcal{E}_0^{(D)}(g), \quad \mathcal{E}_4^{(D)}(g), \quad \mathcal{E}_6^{(D)}(g) \quad : G(\mathbb{Z}) \backslash G(\mathbb{R}) / K \rightarrow \mathbb{C}$$

Automorphic forms

An *automorphic form* is a smooth function $\varphi : G(\mathbb{R}) \rightarrow \mathbb{C}$ satisfying the following conditions

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- (B) φ is an eigenfunction under right-translations of $k \in K$

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- (C) φ is an eigenfunction to all G -invariant differential operators

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- (B) K-finiteness: $\dim(\text{span}\{\varphi(gk) \mid k \in K\}) < \infty$
- (C) Z-finiteness: $\dim(\text{span}\{X\varphi(g) \mid X \in \mathcal{Z}(\mathfrak{g}_{\mathbb{C}})\}) < \infty$

$\mathcal{Z}(\mathfrak{g}_{\mathbb{C}})$ is the center of the universal enveloping algebra $\mathcal{U}(\mathfrak{g}_{\mathbb{C}})$

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- (C) Z-finiteness: $\dim(\text{span}\{X\varphi(g) \mid X \in \mathcal{Z}(\mathfrak{g}_{\mathbb{C}})\}) < \infty$
- (D) Growth: for any norm $\|\cdot\|$ on $G(\mathbb{R})$ there exists a positive integer n and constant C such that $|\varphi(g)| \leq C\|g\|^n$

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Automorphic forms

An *automorphic form* is a smooth function $\varphi : G(\mathbb{R}) \rightarrow \mathbb{C}$ satisfying the following conditions

- (A) Automorphic invariance: ✓ U-duality
- (B) K-finiteness:
- (C) Z-finiteness:
- (D) Growth:

Automorphic forms

An *automorphic form* is a smooth function $\varphi : G(\mathbb{R}) \rightarrow \mathbb{C}$ satisfying the following conditions

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- (B) K-finiteness: ✓ spherical
- (C) Z-finiteness:
- (D) Growth:

Automorphic forms

An *automorphic form* is a smooth function $\varphi : G(\mathbb{R}) \rightarrow \mathbb{C}$ satisfying the following conditions

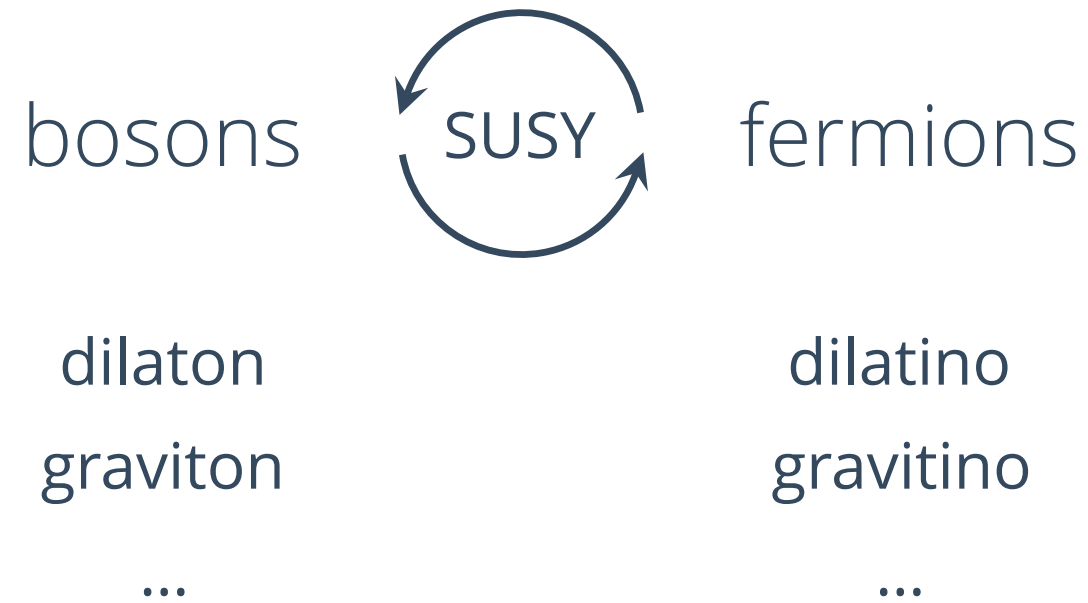
- (A) Automorphic invariance: ✓ U-duality
- (B) K-finiteness: ✓ spherical
- (C) Z-finiteness:
- (D) Growth: ✓ weak coupling limit from string perturbation theory

Automorphic forms

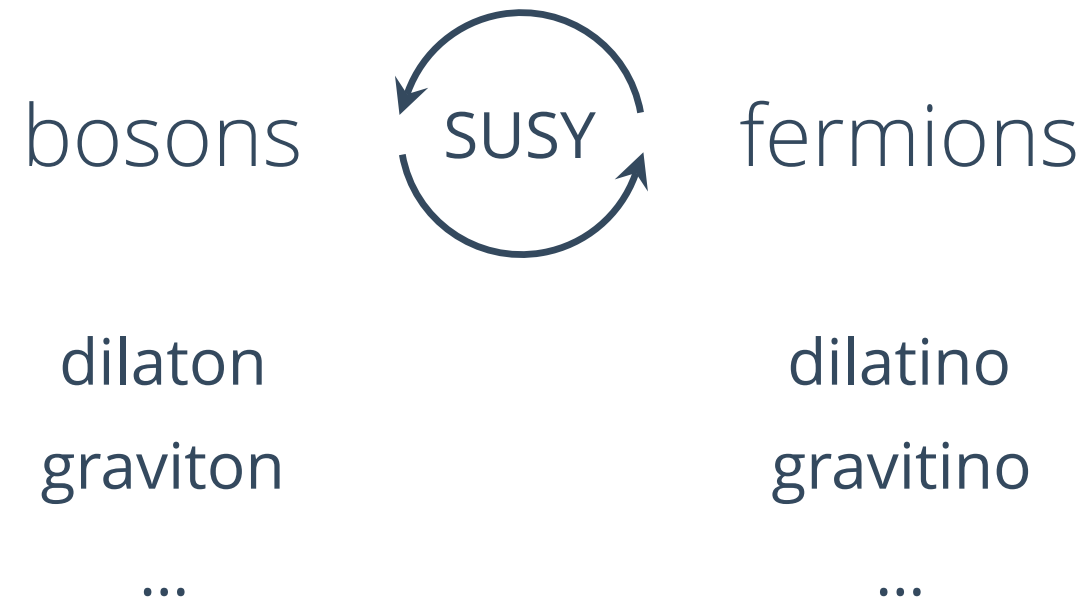
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- (C) Z-finiteness: ?
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Supersymmetry constraints

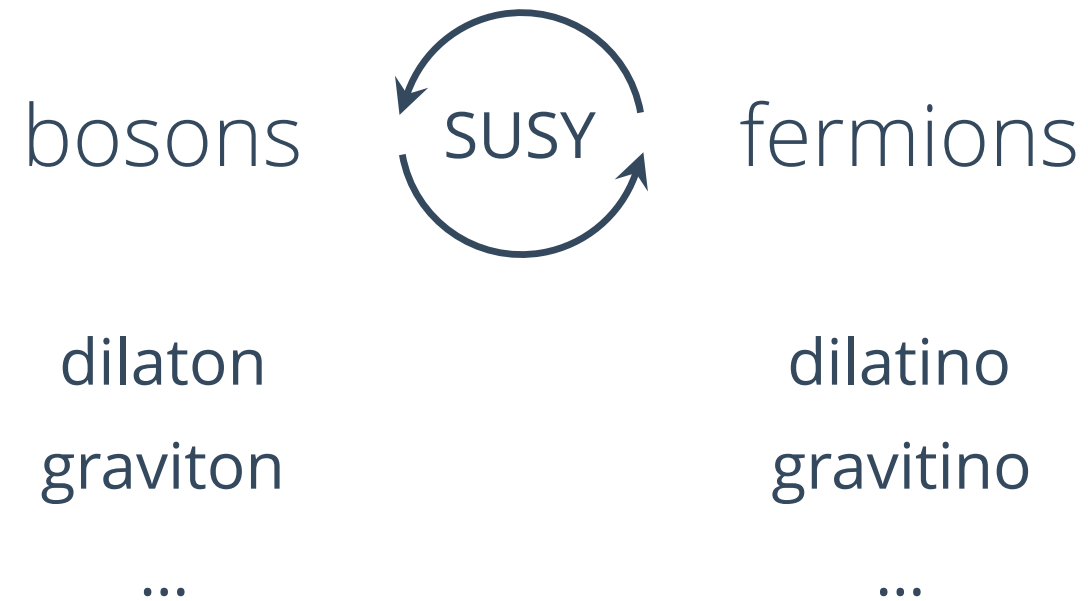


Supersymmetry constraints



10 dimensions:

Supersymmetry constraints

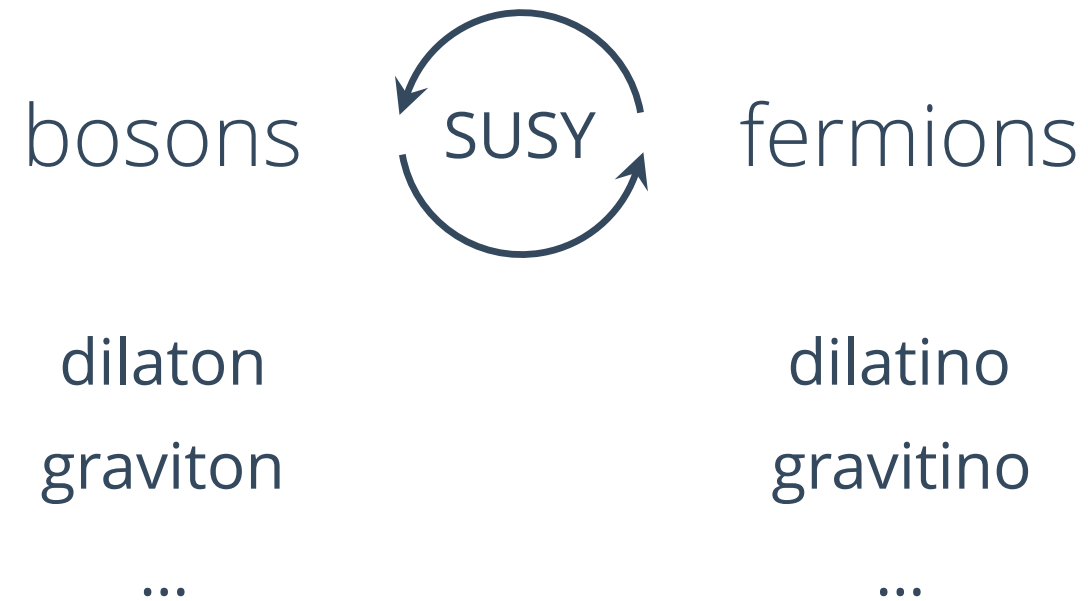


10 dimensions:

$$\mathcal{L}^{(3)} =$$

$$\mathcal{E}_0(\tau) R^4$$

Supersymmetry constraints

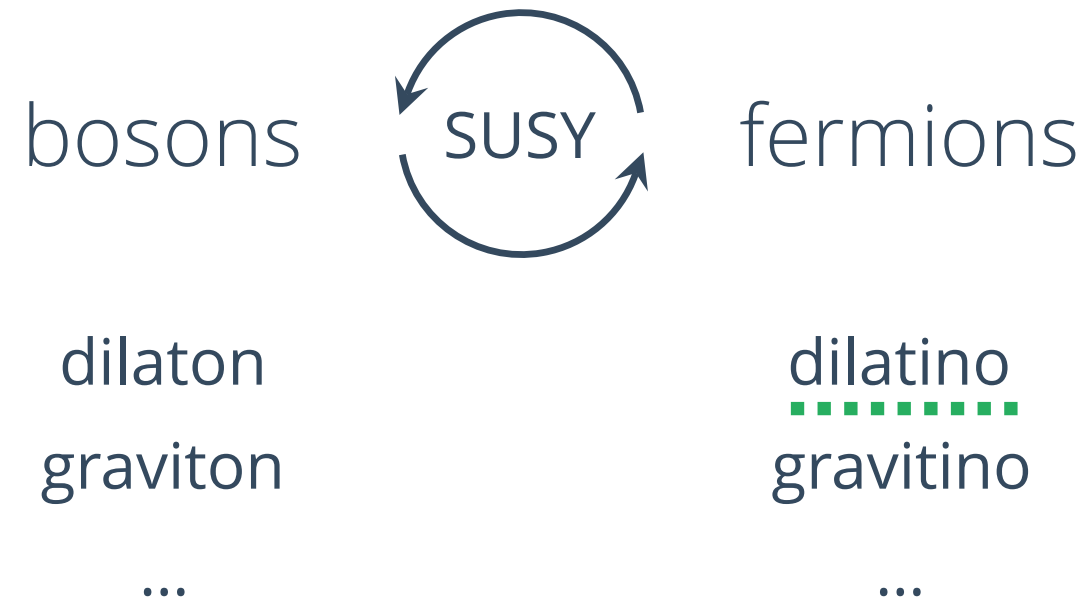


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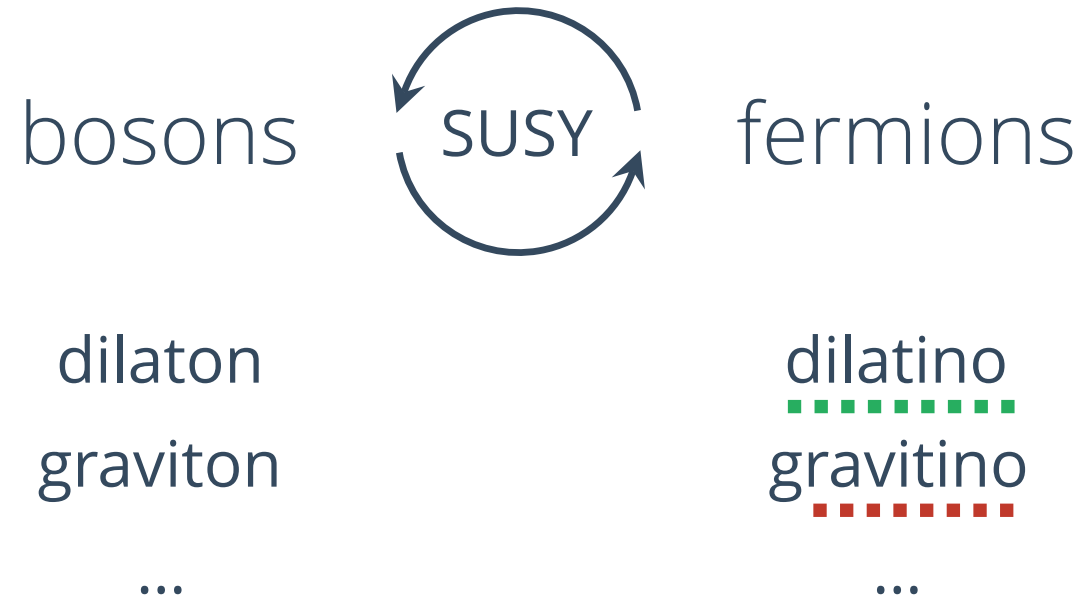


10 dimensions:

$$\mathcal{L}^{(3)} = f_{12}(\tau) \lambda^{\dots 16} + f_{11}(\tau) \hat{G} \lambda^{14} + \dots + f_0(\tau) R^4 + \dots + f_{-12}(\tau) \lambda^{\dots *16}$$

[Green-Sethi]

Supersymmetry constraints

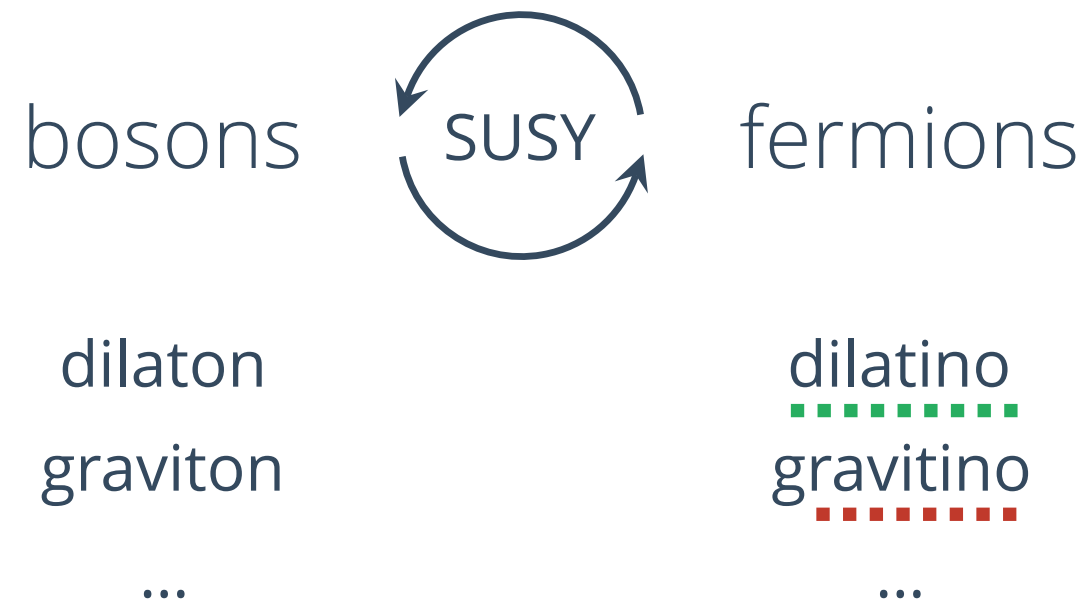


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(Note: In the original image, the λ^{16} term has green dots below it, the $\hat{G}$ term has red dots below it, and the λ^{*16} term has green dots below it.)

Linearized SUSY:

$$f_{w+1}(\tau) = i \left(\tau_2 \frac{\partial}{\partial \tau} - i \frac{w}{2} \right) f_w(\tau)$$

[Green-Sethi]

Supersymmetry constraints

$$\int d^D x \sqrt{-g} \mathcal{L} = S = S^{(0)}$$
$$\delta \Psi = \delta^{(0)} \Psi$$

Supersymmetry constraints

$$\int d^D x \sqrt{-g} \mathcal{L} = S = S^{(0)} + (\alpha')^3 S^{(3)} + (\alpha')^5 S^{(5)} + \dots$$
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$$\begin{cases} \delta^{(0)} S^{(3)} + \delta^{(3)} S^{(0)} = 0 \\ [\delta_1, \delta_2] \lambda^* = \delta_{\text{local symmetries}} \lambda^* + (\text{equations of motion}) \end{cases}$$

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$$\implies (\Delta - \frac{3}{4}) \mathcal{E}_0(\tau) = 0$$

[Green-Sethi]

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[Green-Sethi]

Supersymmetry constraints



10 dimensions: $\Delta = 4\tau_2^2 \frac{\partial}{\partial \tau} \frac{\partial}{\partial \bar{\tau}}$ Laplacian on
Poincaré UHP

$$(\Delta - \frac{3}{4})\mathcal{E}_0(\tau) = 0$$

[Green-Sethi]

Supersymmetry constraints



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$$(\Delta - 12)\mathcal{E}_6(\tau) = -(\mathcal{E}_0(\tau))^2 \quad \text{[Green-Vanhove]}$$

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Not an automorphic form in a strict sense

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Similarly for lower dimensions

Eisenstein series

$$E(s; \tau) =$$

$$s \in \mathbb{C}$$

Eisenstein series

$$E(s; \tau) = \sum_{\substack{c, d \in \mathbb{Z} \\ (c, d) \neq (0, 0)}} \frac{\operatorname{Im}(\tau)^s}{|c\tau + d|^{2s}}$$

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$$B(\mathbb{Z}) = \left\{ \begin{pmatrix} * & * \\ 0 & * \end{pmatrix} \in SL(2, \mathbb{Z}) \right\} \quad \text{Borel subgroup}$$

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$$\tau \mapsto \operatorname{Im}(\tau)^s$$

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Completed Riemann zeta function

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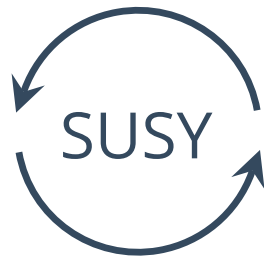
Eisenstein series

$$(\Delta - s(s-1))E(s; \tau) = 0 \qquad E(s; \tau) \sim \tau_2^s \qquad g_s = \tau_2^{-1} \rightarrow 0$$

[Green-Gutperle, Pioline, Green-Russo-Vanhove]

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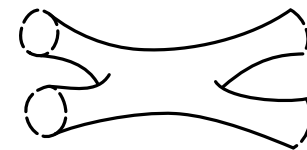
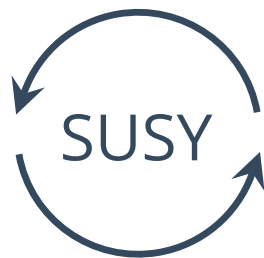
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[Green-Gutperle, Pioline, Green-Russo-Vanhove]

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(Einstein frame)

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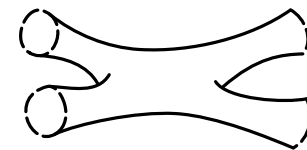
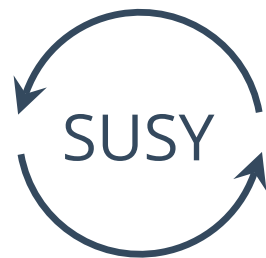
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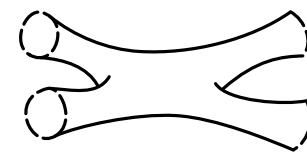
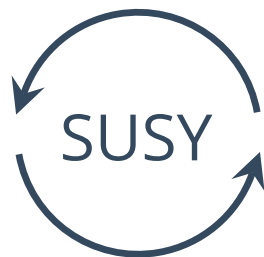
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$\mathcal{E}_6(\tau)$ as a sum over images $\sum_{B(\mathbb{Q}) \backslash G(\mathbb{Z})}$ but not of a character χ

[Green-Miller-Vanhove]

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Expand Bessel function in g_s

$$\tau = \chi + i g_s^{-1}$$

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Perturbative
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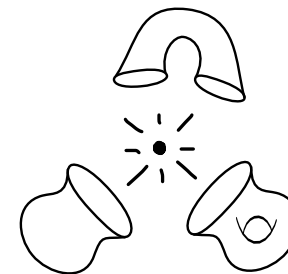
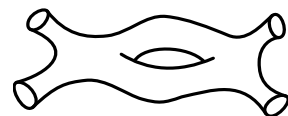
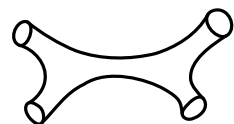
Non-perturbative
(remaining modes)

[Green-Gutperle]

Extracting physical information

Expand Bessel function in g_s

$$\tau = \chi + i g_s^{-1}$$



$$\mathcal{E}_0(\tau) = 2\zeta(3)g_s^{-3/2} + 4\zeta(2)g_s^{1/2} + 2\pi \sum_{m \neq 0} \sqrt{|m|} \sigma_{-2}(m) e^{-2\pi|m|g_s^{-1} + 2\pi i m \chi} \left[1 + \mathcal{O}(g_s) \right]$$

Perturbative
(zero-mode)

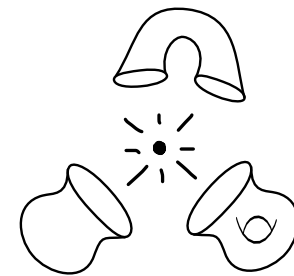
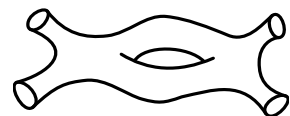
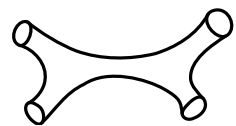
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Instanton action

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Perturbative
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[Green-Gutperle]

Extracting physical information

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.....

Perturbative
(zero-mode)

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Non-perturbative
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$\sigma_s(m) = \sum_{d|m} d^s$

Sums over the number of ways the charge m can be factorised into two integers

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Sums over the number of ways the charge m can be factorised into two integers



wrapping number and charge
of a T-dual D-particle

[Green-Gutperle]

Lower dimensions

Lower dimensions

D	$G(\mathbb{R})$	K	$G(\mathbb{Z})$
10	$SL(2, \mathbb{R})$	$SO(2)$	$SL(2, \mathbb{Z})$
9	$SL(2, \mathbb{R}) \times \mathbb{R}^+$	$SO(2)$	$SL(2, \mathbb{Z}) \times \mathbb{Z}_2$
8	$SL(3, \mathbb{R}) \times SL(2, \mathbb{R})$	$SO(3) \times SO(2)$	$SL(3, \mathbb{Z}) \times SL(2, \mathbb{Z})$
7	$SL(5, \mathbb{R})$	$SO(5)$	$SL(5, \mathbb{Z})$
6	$Spin(5, 5; \mathbb{R})$	$(Spin(5) \times Spin(5)) / \mathbb{Z}_2$	$Spin(5, 5; \mathbb{Z})$
5	$E_6(\mathbb{R})$	$USp(8) / \mathbb{Z}_2$	$E_6(\mathbb{Z})$
4	$E_7(\mathbb{R})$	$SU(8) / \mathbb{Z}_2$	$E_7(\mathbb{Z})$
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3	$E_8(\mathbb{R})$	$Spin(16) / \mathbb{Z}_2$	$E_8(\mathbb{Z})$

$$E(\chi; g) = \sum_{\gamma \in B(\mathbb{Z}) \setminus G(\mathbb{Z})} \chi(\gamma g)$$

Parabolic subgroups

Fourier expand
in different directions



Unipotent subgroup U

Parabolic subgroups

Fourier expand
in different directions



Unipotent subgroup U



Choice of parabolic subgroup P

Parabolic subgroups

Fourier expand
in different directions



Unipotent subgroup U



Choice of parabolic subgroup P

Σ choice of simple roots

$\langle \Sigma \rangle$ generated root system

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$$G = SL(4)$$



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$$P = LU = \left\{ \begin{pmatrix} * & * & * & * \\ * & * & * & * \\ & * & * & * \\ & & * & * \end{pmatrix} \right\} \quad L = \left\{ \begin{pmatrix} * & * & & \\ * & * & & \\ & * & * & \\ & & * & * \end{pmatrix} \right\} \quad U = \left\{ \begin{pmatrix} 1 & * & * & \\ & 1 & * & * \\ & & 1 & * \\ & & & 1 \end{pmatrix} \right\}$$

Parabolic subgroups



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Parabolic subgroups



Parabolic subgroups



Maximal parabolic

Parabolic subgroups



Minimal parabolic
Borel



Maximal parabolic

Parabolic subgroups



Minimal parabolic
Borel

$$B = NA$$

$$N = \left\{ \begin{pmatrix} \boxed{1} & * & * & * \\ & \boxed{1} & * & * \\ & & \boxed{1} & * \\ & & & \boxed{1} \end{pmatrix} \right\}$$



Maximal parabolic

$$P = LU$$

$$U = \left\{ \begin{pmatrix} \boxed{1} & & & * \\ & \boxed{1} & & * \\ & & \boxed{1} & * \\ & & & \boxed{1} \end{pmatrix} \right\}$$

Fourier expansion

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Let $\psi : U(\mathbb{Z}) \backslash U(\mathbb{R}) \rightarrow U(1)$ be a multiplicative character

Parametrised by $m_\alpha \in \mathbb{Z}$ called charges

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Fourier expansion

Let $\psi : U(\mathbb{Z}) \backslash U(\mathbb{R}) \rightarrow U(1) \cong S^1$ be a multiplicative character

Parametrised by $m_\alpha \in \mathbb{Z}$ called charges



$$\psi_U \left(\begin{pmatrix} 1 & & & y_1 \\ & 1 & & y_2 \\ & & 1 & y_3 \\ & & & 1 \end{pmatrix} \right) = e^{2\pi i(m_1 y_1 + m_2 y_2 + m_3 y_3)}$$

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$\cong S^1$

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$$F_U(\chi, \psi; g) = \int_{U(\mathbb{Z}) \backslash U(\mathbb{R})} E(\chi, ug) \overline{\psi(u)} du$$

Fourier expansion

Fourier expansion

$$E(\chi; g) = \sum_{\psi} F_U(\chi, \psi; g)$$

Fourier expansion

$$E(\chi; g) = F_U(\chi, 1; g) + \sum_{\psi \neq 1} F_U(\chi, \psi; g)$$

Fourier expansion

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Fourier expansion

$$E(\chi; g) = F_U(\chi, 1; g) + \sum_{\psi^{(1)} \neq 1} F_{U^{(1)}}(\chi, \psi^{(1)}; g) + \sum_{\psi^{(2)} \neq 1} F_{U^{(2)}}(\chi, \psi^{(2)}; g) + \dots$$

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Terminology

$P = B \longrightarrow U = N$ Fourier coefficient is a Whittaker coefficient

$$N = \left\{ \begin{pmatrix} 1 & * & * & * \\ & 1 & * & * \\ & & 1 & * \\ & & & 1 \end{pmatrix} \right\}$$

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$$F_U$$

$$W_N$$

Terminology

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Characters and coefficients with all $m_\alpha \neq 0$ are called **generic**
otherwise they are called **degenerate**

Fourier expansion

Choice of unipotent subgroup U $\longleftrightarrow$ Study different perturbative and non-perturbative effects

[Green-Miller-Vanhove]

Fourier expansion

Choice of unipotent subgroup U $\longleftrightarrow$ Study different perturbative and non-perturbative effects

- String perturbation limit
D-instantons | NS5-instantons

$$g_s \rightarrow 0$$



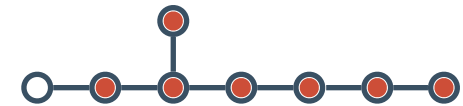
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- Decompactification limit
Higher dimensional black holes | BPS states

Large radius for
compactified circle



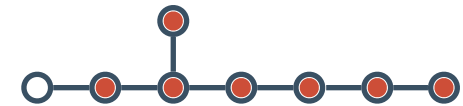
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- M-theory limit
M2, M5-instantons

Large M-theory torus



[Green-Miller-Vanhove]

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[Green-Miller-Vanhove]

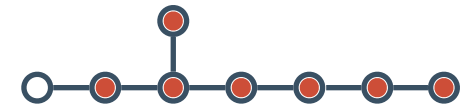
Maximal parabolic
subgroups

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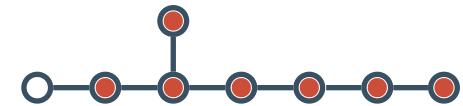
Difficult to compute!

Fourier expansion

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[Green-Miller-Vanhove]

Maximal parabolic
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Difficult to compute!

Recent result in [Bossard-Pioline]

Fourier expansion

Goal: find expressions for Fourier coefficients
in terms of (known) Whittaker coefficients

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Goal: find expressions for Fourier coefficients
in terms of (known) Whittaker coefficients

Would allow us to compute non-perturbative effects that
capture information about instantons and black holes

Adelic framework

*An **efficient**, but abstract, way to approach the subject of automorphic forms is by the introduction of **adeles**, rather **ungainly objects** that nevertheless, once familiar, **spare** much unnecessary thought and **many useless calculations**.*

— Robert P. Langlands*

*Representation theory - its rise and its role in number theory, Proceedings of the Gibbs symposium (1989)

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Eisenstein series

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— Robert P. Langlands*

Adelic Eisenstein series



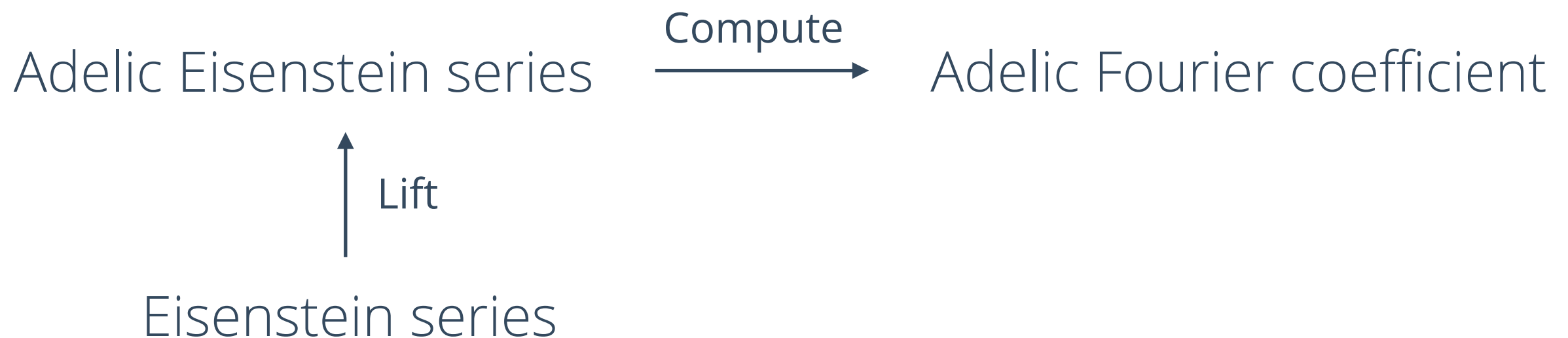
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The adeles

For a prime p

$\mathbb{Q}$

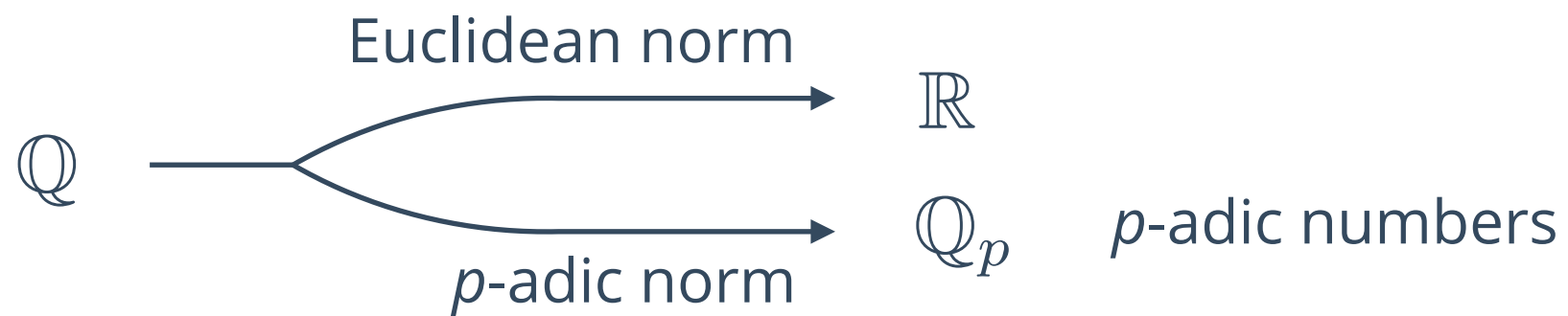
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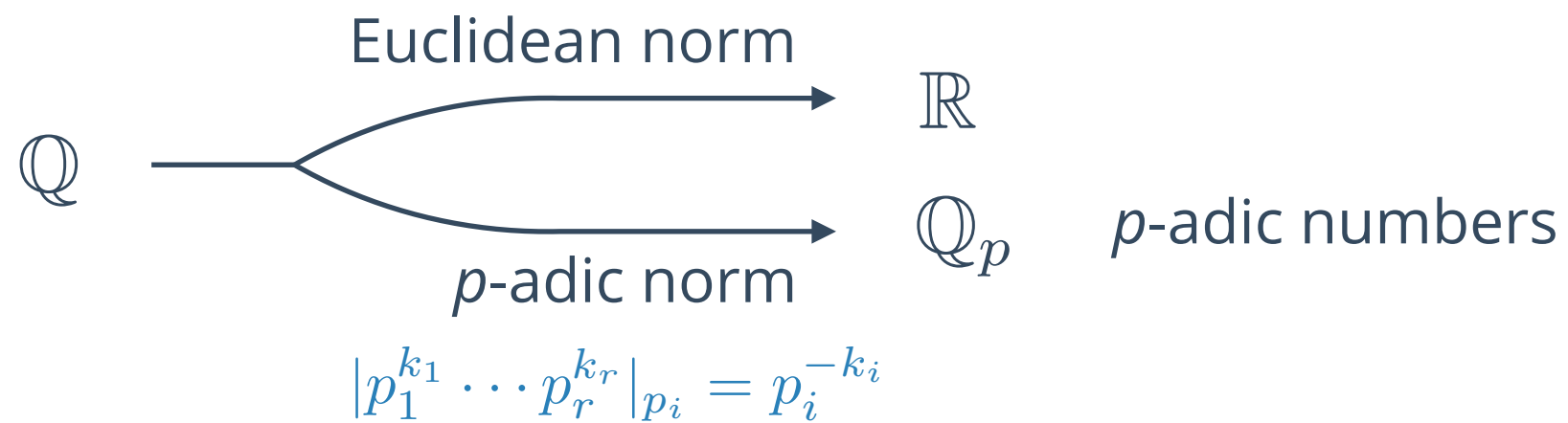
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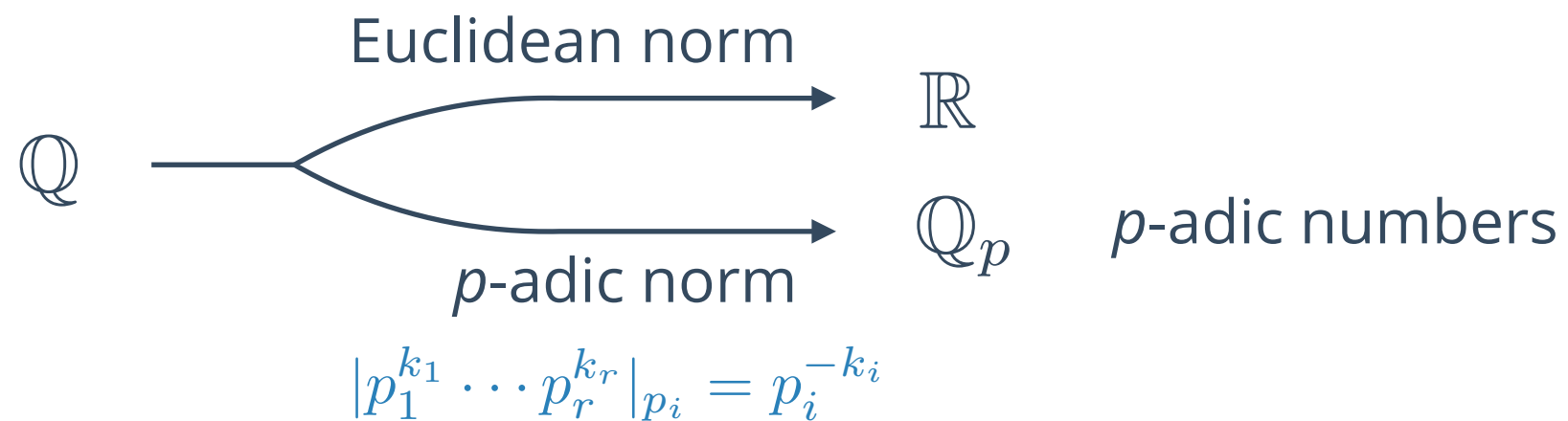
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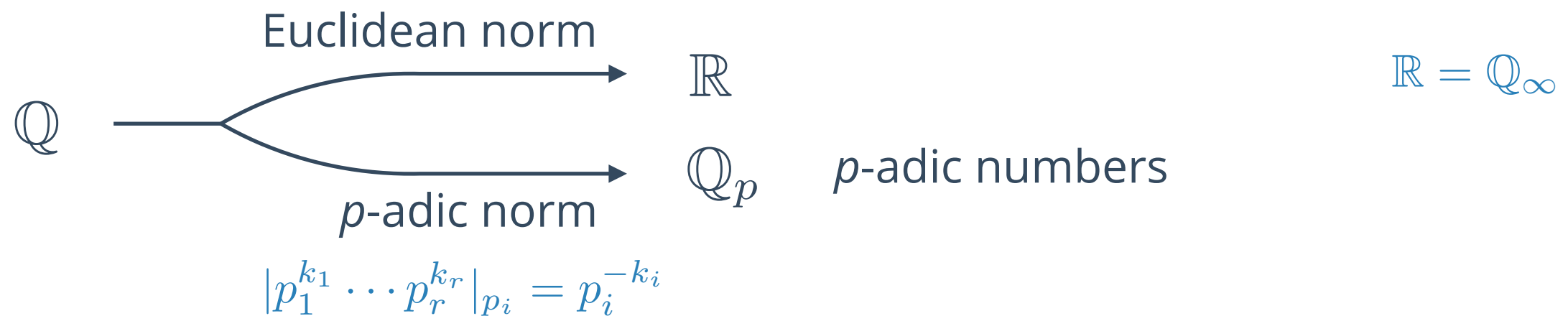
For a prime p



$$\mathbb{R} = \mathbb{Q}_\infty$$

The adeles

For a prime p

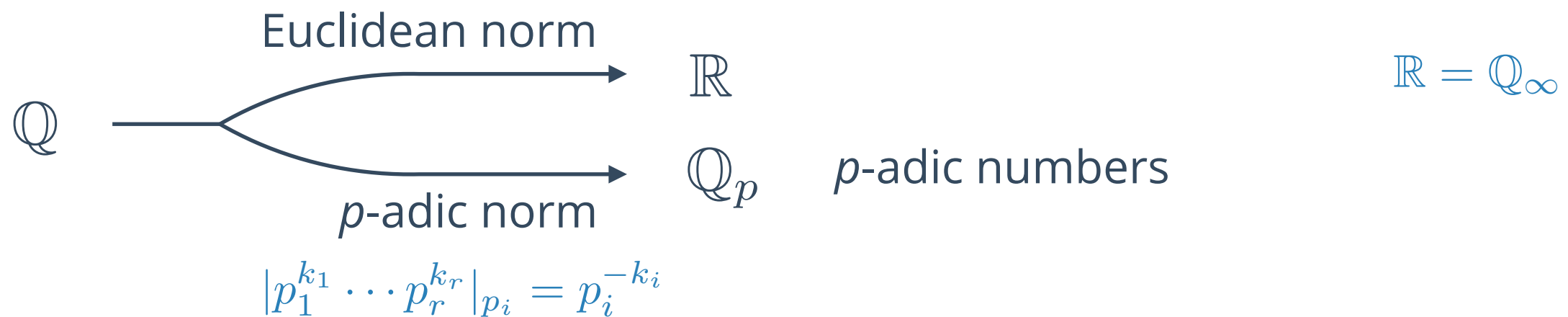


The adeles are then defined as

$$\mathbb{A} = \mathbb{A}_{\mathbb{Q}} = \mathbb{R} \times \prod'_{p \text{ prime}} \mathbb{Q}_p \quad x = (x_\infty; x_2, x_3, x_5, \dots) \in \mathbb{A}$$

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$$\mathbb{Q} \hookrightarrow \mathbb{A}$$

$$q \mapsto (q; q, q, \dots)$$

$\mathbb{Q}$ is discrete in $\mathbb{A}$ taking the role of $\mathbb{Z}$ in $\mathbb{R}$

Much easier to work with since it is a field!

Adelic framework

$$\mathcal{E}_0^{(D)}(g), \quad \mathcal{E}_4^{(D)}(g), \quad \mathcal{E}_6^{(D)}(g) \quad : G(\mathbb{Z}) \backslash G(\mathbb{R}) / K \rightarrow \mathbb{C}$$

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Lift to the adeles

[FGKP15 §4.2.2]

$$G(\mathbb{A}) = G(\mathbb{R}) \times \prod'_{p \text{ prime}} G(\mathbb{Q}_p) \quad K_{\mathbb{A}} = K \times \prod_{p \text{ prime}} G(\mathbb{Z}_p)$$

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Eisenstein series $\longrightarrow$ Adelic Eisenstein series

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$$\int_{U(\mathbb{Z}) \backslash U(\mathbb{R})} E(\chi; ug) \overline{\psi_{\mathbb{R}}(u)} du$$

$$\int_{U(\mathbb{Q}) \backslash U(\mathbb{A})} E(\chi; ug) \overline{\psi_{\mathbb{A}}(u)} du$$

$$m_{\alpha} \in \mathbb{Z}$$

$$m_{\alpha} \in \mathbb{Q}$$

Computing adelic Fourier coefficients

[FGKP15 §9-10]

Whittaker coefficients

Computing adelic Fourier coefficients

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Constant term: Langlands' constant term formula

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Maximally degenerate

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Factorisation

[GKP14]

Fourier coefficients

In terms of Whittaker coefficients

Simplify drastically for certain χ

Example of simplifications

$$G = SL(3)$$

$$E(\chi; g)$$

$$\chi \longleftrightarrow (s_1, s_2) \in \mathbb{C}^2$$

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$$N = \left\{ \begin{pmatrix} 1 & * & * \\ & 1 & * \\ & & 1 \end{pmatrix} \right\}$$

$$\psi_{m_1, m_2} \left(\begin{pmatrix} 1 & x_1 & * \\ & 1 & x_2 \\ & & 1 \end{pmatrix} \right) = e^{2\pi i(m_1 x_1 + m_2 x_2)}$$

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$$W_N(\chi, \psi_{m_1, m_2}; g) \propto \left(\text{arithmetic factor} \right) \int K_{\#}(\dots) K_{\#}(\dots)$$

[FGKP15 §10.6]

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└ p-adic part
↓

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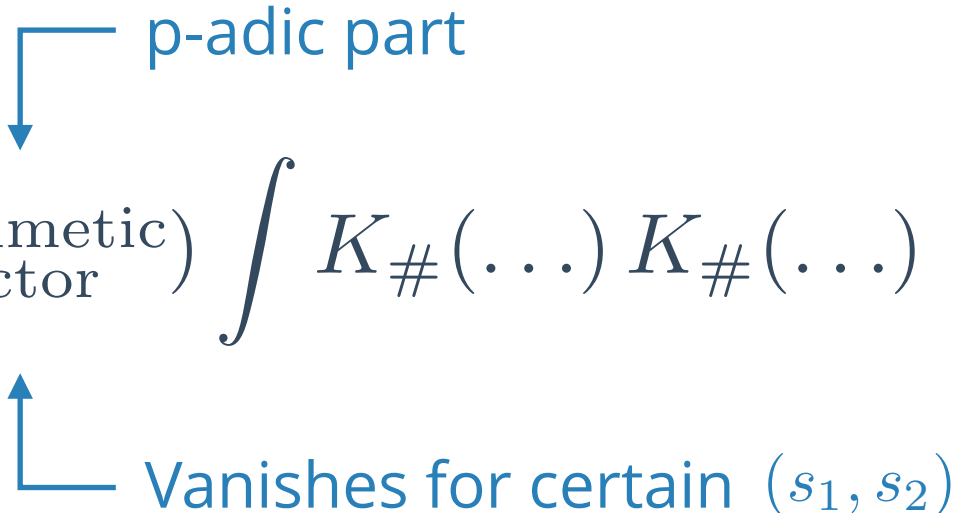
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To explain this, we need to study
small automorphic representations

[FGKP15 §10.6]

Automorphic representations

$G(\mathbb{A}) \curvearrowright$ Space of automorphic forms*

* With some subtleties described in [FGKP15 §6]

[Bump, Goldfeld-Hundley]

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What is a small automorphic representation?

* With some subtleties described in [FGKP15 §6]

[Bump, Goldfeld-Hundley]

Wavefront set

[Mœglin–Waldspurger, Matumoto, Ginzburg-Rallis-Soudry, Ginzburg,
Gomez-Gourevitch-Sahì, Jiang-Liu-Savin, Joseph, Miller-Sahì]

Wavefront set

The (global) wavefront set contains all the characters ψ which can give rise to non-vanishing Fourier coefficients in that representation

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Small automorphic representations have
few non-vanishing Fourier coefficients

[Mœglin–Waldspurger, Matumoto, Ginzburg-Rallis-Soudry, Ginzburg,
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Characters $\psi \longleftrightarrow$ Nilpotent elements in $\mathfrak{g}(\mathbb{Q})$

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Closure with respect to partial ordering

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[Collingwood-McGovern]

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Illustrated by a Hasse diagram

Nilpotent orbits

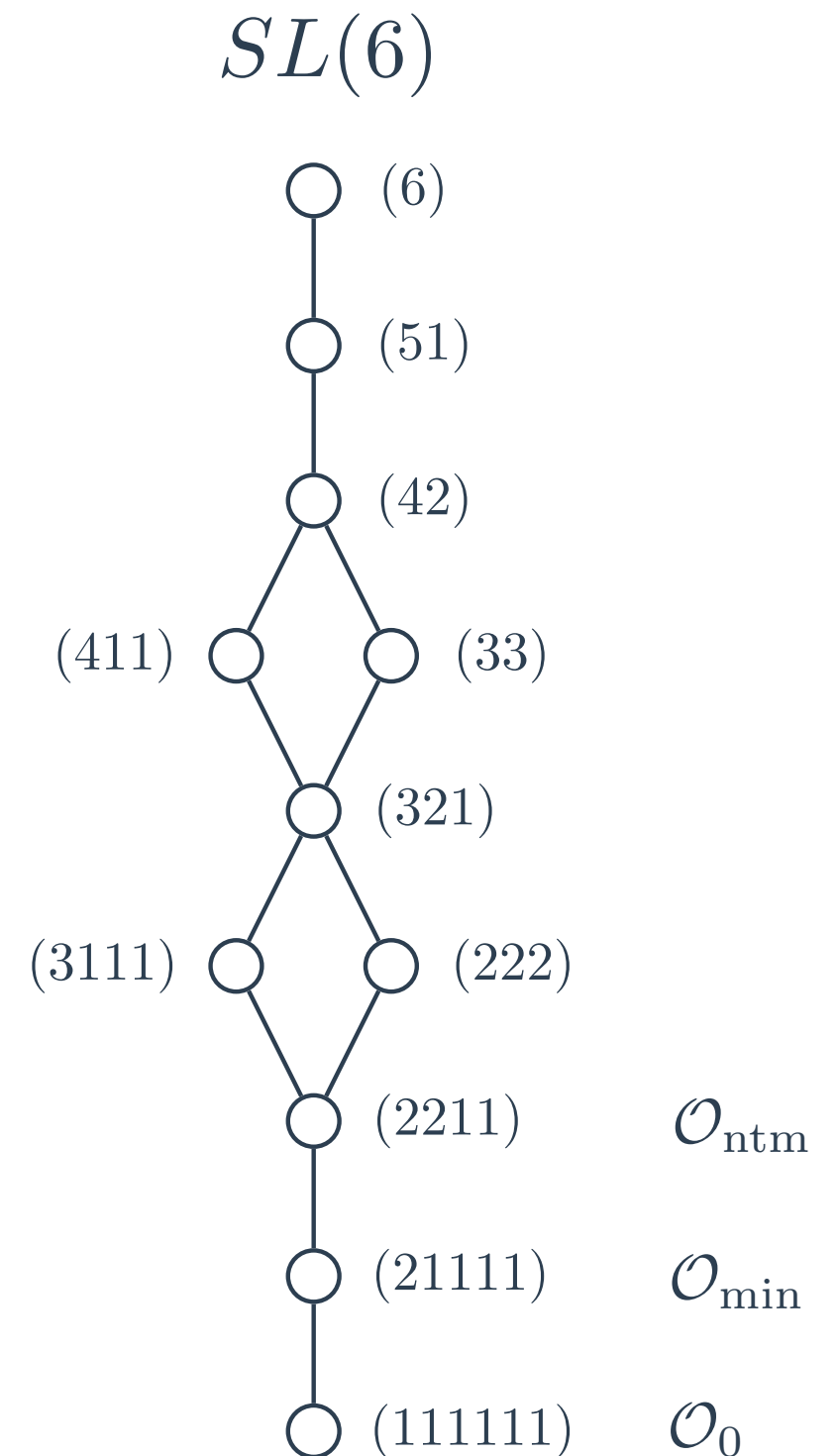
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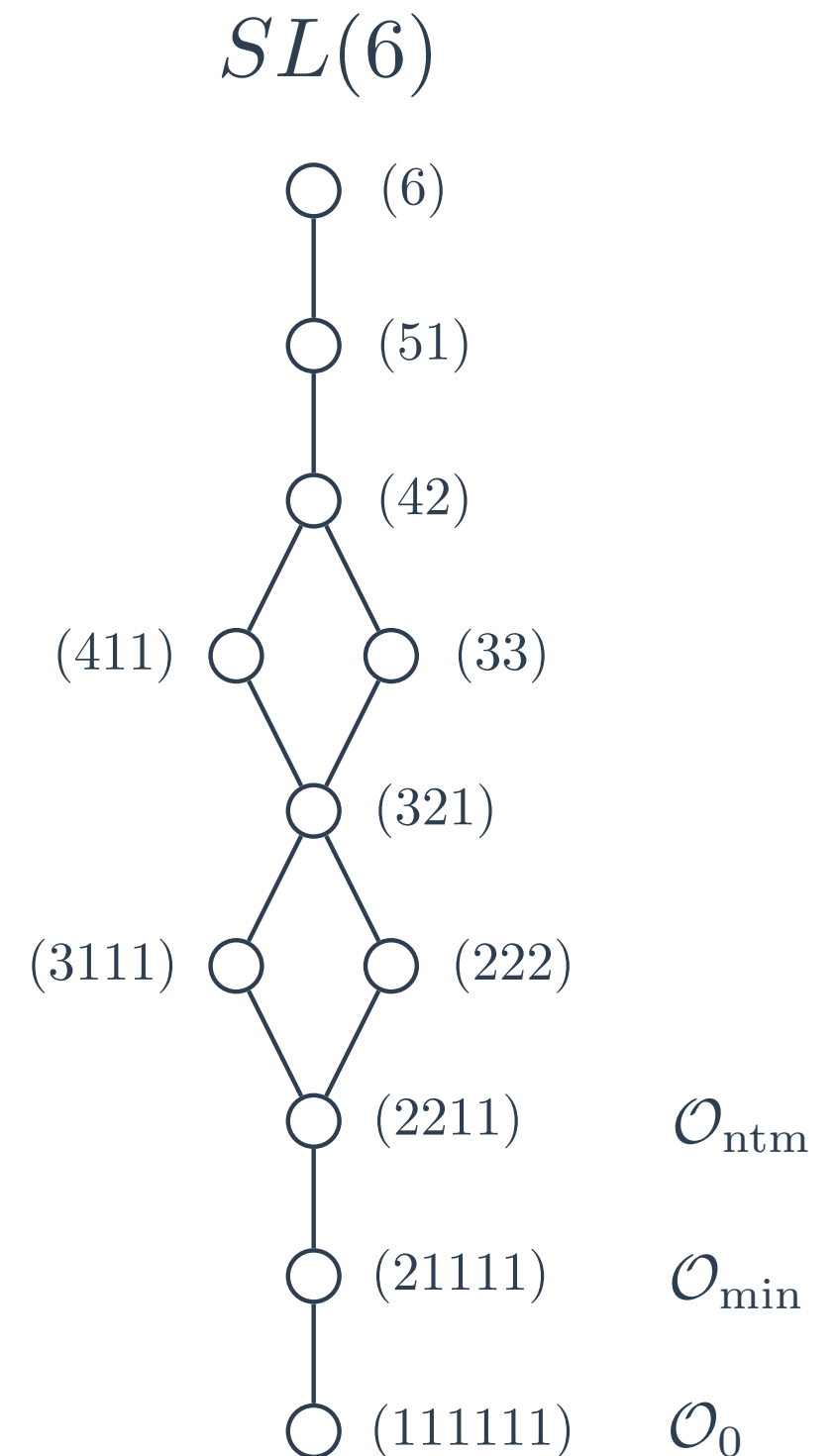
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Illustrated by a Hasse diagram

Closure: $\overline{\mathcal{O}} = \bigcup_{\mathcal{O}' \leq \mathcal{O}} \mathcal{O}'$



Automorphic representations

Small representations

Automorphic representations

Small representations

$$\mathrm{WF}(\pi_{\min}) = \overline{\mathcal{O}_{\min}} = \mathcal{O}_{\min} \cup \mathcal{O}_0$$

$$\mathrm{WF}(\pi_{\mathrm{ntm}}) = \overline{\mathcal{O}_{\mathrm{ntm}}} = \mathcal{O}_{\mathrm{ntm}} \cup \mathcal{O}_{\min} \cup \mathcal{O}_0$$

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[Green-Miller-Vanhove,
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$$\int K K \longrightarrow 0$$

$$\sum K \longrightarrow K$$

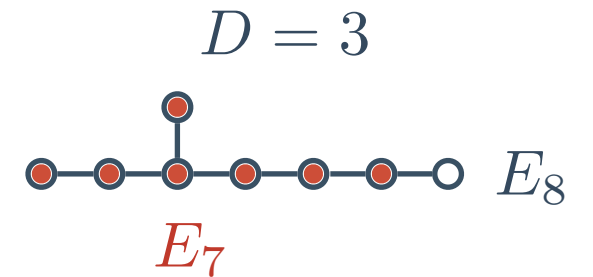
Automorphic representations

[Ferrara-Günaydin, Ferrara-Maldacena, Green-Miller-Vanhove]

Automorphic representations

- Decompactification limit
Higher dimensional black holes | BPS states

Large radius for
compactified circle

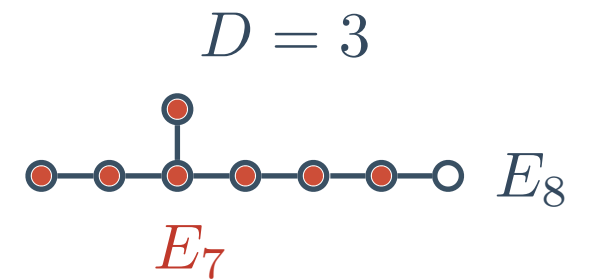


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$\pi_{\min}$

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π_{3A_1}

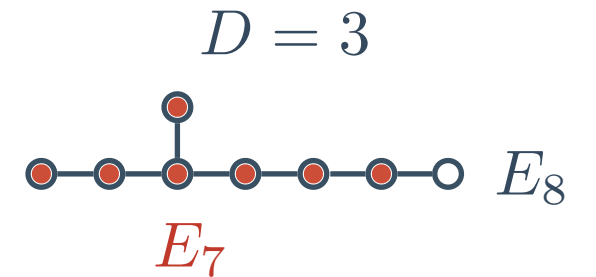
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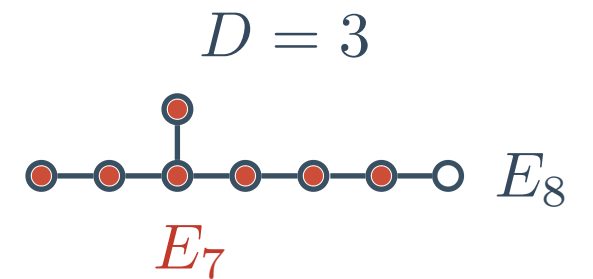
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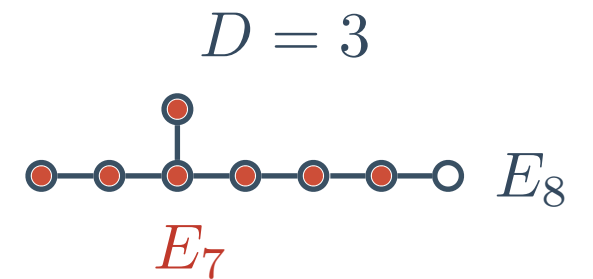
[Miller-Sahi]

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$\dim\{\psi_U \in \text{WF}\}$	28	45	55	56
	[Miller-Sahi]			
$D = 4$				
BPS-orbits	$\frac{1}{2}$ BPS	$\frac{1}{4}$ BPS	$\frac{1}{8}$ BPS	$\frac{1}{8}$ BPS ⁺
$E_7 \curvearrowright \{\psi_U\}$				

[Ferrara-Günaydin, Ferrara-Maldacena, Green-Miller-Vanhove]

Goal: find expressions for Fourier coefficients
in terms of (known) Whittaker coefficients
using vanishing properties of the given π

Previous results

[Miller-Sahai]

Previous results

Theorem

For $G = E_6, E_7$, an automorphic form $\varphi \in \pi_{\min}$ is completely determined by maximally degenerate Whittaker coefficients

$$W_N \quad \text{with only one } m_\alpha \neq 0$$

[Miller-Sahi]

Main results

$SL(3), SL(4)$

[GKP14]

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$SL(3), SL(4)$

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More generally, for $\varphi \in \pi$

[GKP14]

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$$\varphi = \sum_{\mathcal{O}} \varphi_{\mathcal{O}} \quad \text{where } \varphi_{\mathcal{O}} \text{ vanishes unless } \mathcal{O} \subseteq \text{WF}(\pi)$$

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Main results *SL(3), SL(4)*

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Corollary

[GKP14]

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Corollary

$\varphi \in \pi_{\min}$ maximally degenerate Whittaker coefficients

[GKP14]

Main results $SL(3), SL(4)$

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Corollary

$\varphi \in \pi_{\min}$ maximally degenerate Whittaker coefficients single root

[GKP14]

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single root

[GKP14]

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$$\varphi = \sum_{\mathcal{O}} \varphi_{\mathcal{O}} \quad \text{where } \varphi_{\mathcal{O}} \text{ vanishes unless } \mathcal{O} \subseteq \text{WF}(\pi)$$

Corollary

$$\varphi \in \pi_{\min}$$

single root

$$\varphi \in \pi_{\text{ntm}}$$

at most two commuting roots

[GKP14]

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$SL(3), SL(4)$

Fourier coefficients on maximal parabolic subgroups in the minimal representation



$\pi_{\min}$

[GKP14]

Main results

$SL(3), SL(4)$

Fourier coefficients on maximal parabolic subgroups in the minimal representation



$\pi_{\min}$

Theorem

$$F_U(\chi_{\min}, \psi; g) = W_N(\chi_{\min}, \psi'; lg) \quad \text{with } l \in L(\mathbb{Q}) \text{ depending on } \psi$$

[GKP14]

Main results

$SL(3), SL(4)$

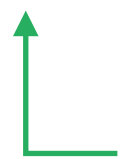
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Maximal parabolic
Fourier coefficient

[GKP14]

Main results

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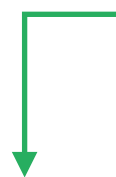
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Known Whittaker coefficient



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Maximal parabolic
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[GKP14]

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Maximal parabolic
Fourier coefficient

Maximally degenerate

[GKP14]

Other groups

[Work in progress with Ahlén, Liu, Kleinschmidt, Persson]

Other groups

$$SL(n)$$

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Maximal parabolic
Fourier coefficient



Maximally degenerate

and similar statement for next-to-minimal representation

[Work in progress with Ahlén, Liu, Kleinschmidt, Persson]

Other groups

Conjecture

A similar relations holds for all simple Lie groups

$$F_U(\chi_{\min}, \psi; g) = W_N(\chi_{\min}, \psi'; lg) \quad \text{with } l \in L(\mathbb{Q}) \text{ depending on } \psi$$

↑
Maximal parabolic
Fourier coefficient

↑
Maximally degenerate

[GKP14]

[Proof in progress with Gourevitch, Kleinschmidt, Persson, Sahi]

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Would allow us to compute non-perturbative effects that capture information about instantons and black holes

[GKP14]

[Proof in progress with Gourevitch, Kleinschmidt, Persson, Sahi]

Outlook

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- ◎ Other compactifications leading to automorphic forms on other groups.

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How to define “small automorphic representations” for Kac-Moody groups? What is the mechanism behind the vanishing properties?

- $\mathcal{E}_6 D^6 R^4$ requires extended notion of automorphic forms, the development of which will positively bring new exciting insights to both physics and mathematics.

Thank you!

Henrik Gustafsson